

Handwritten text in a stylized, cursive script, possibly representing a name or a title. The characters are dark and feature decorative elements like yellow and blue highlights.

Forward-Backward  
Multiplicity Correlations  
in Au+Au

Peter Steinberg for the PHOBOS Collaboration

Chemistry Department  
Brookhaven National Laboratory

Quark Matter 2005, Budapest, Hungary

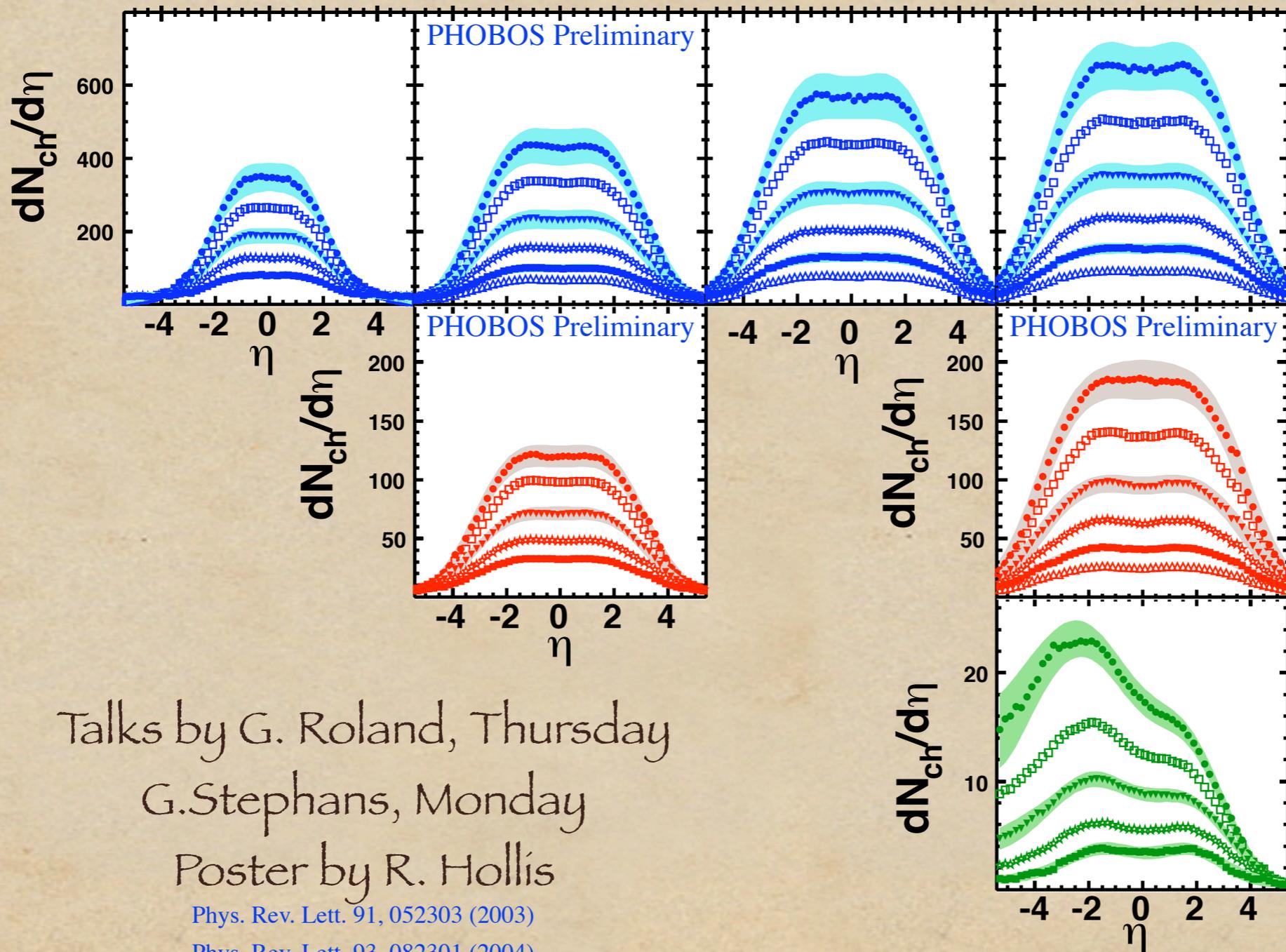
# PHOBOS Collaboration

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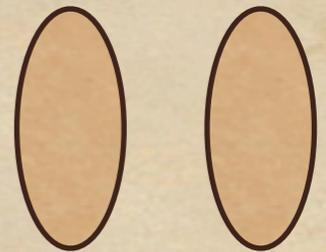


# Single-Particle $dN/d\eta$

19.6 GeV    62.4 GeV    130 GeV    200 GeV



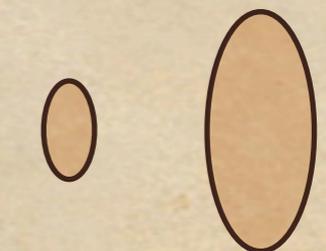
Au+Au



Cu+Cu



d+Au



Talks by G. Roland, Thursday

G. Stephans, Monday

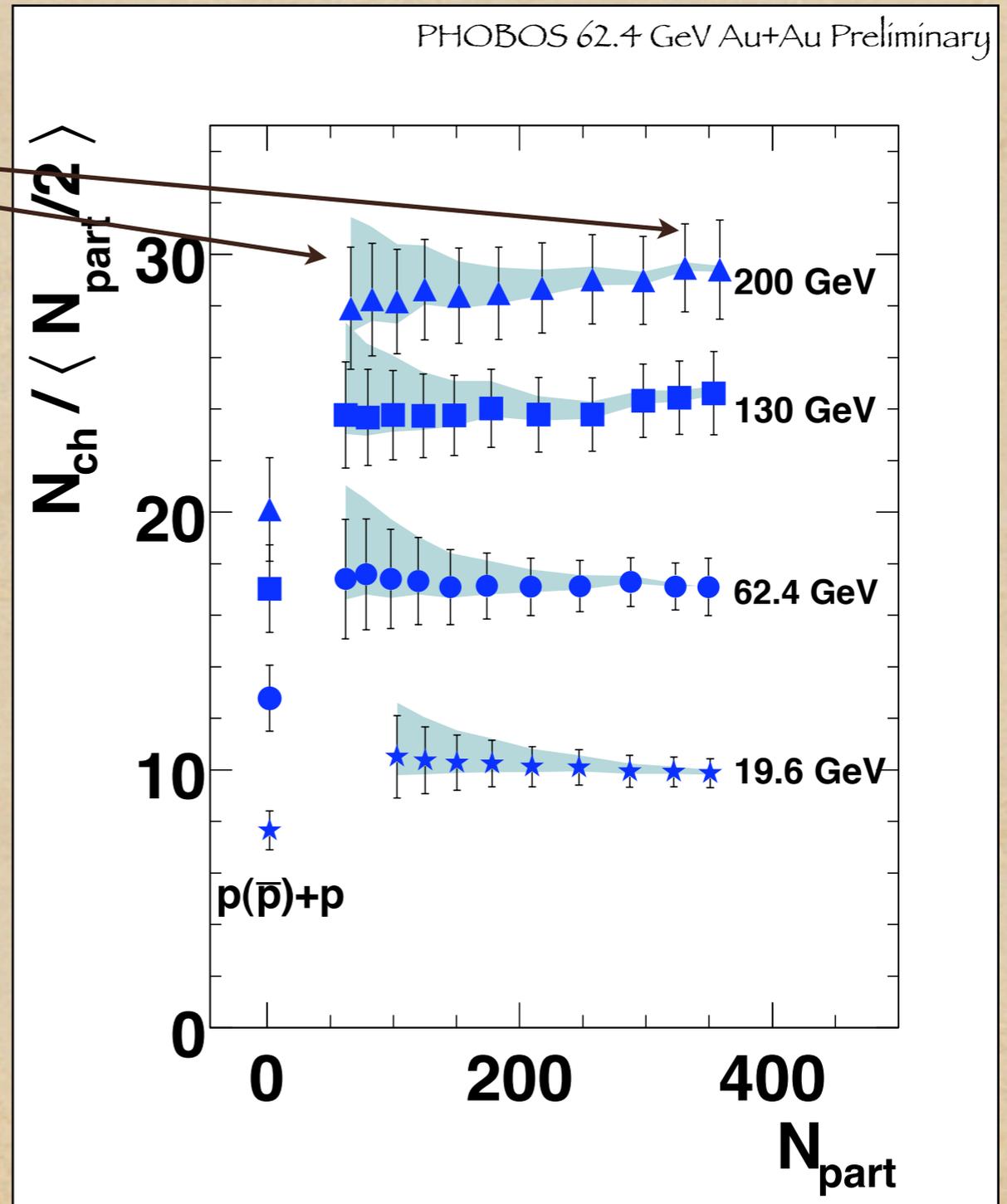
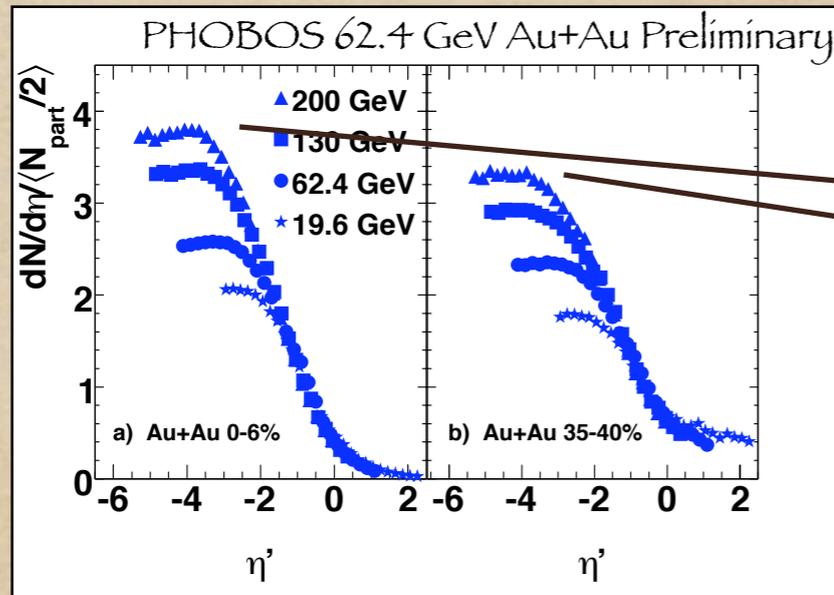
Poster by R. Hollis

[Phys. Rev. Lett. 91, 052303 \(2003\)](#)

[Phys. Rev. Lett. 93, 082301 \(2004\)](#)

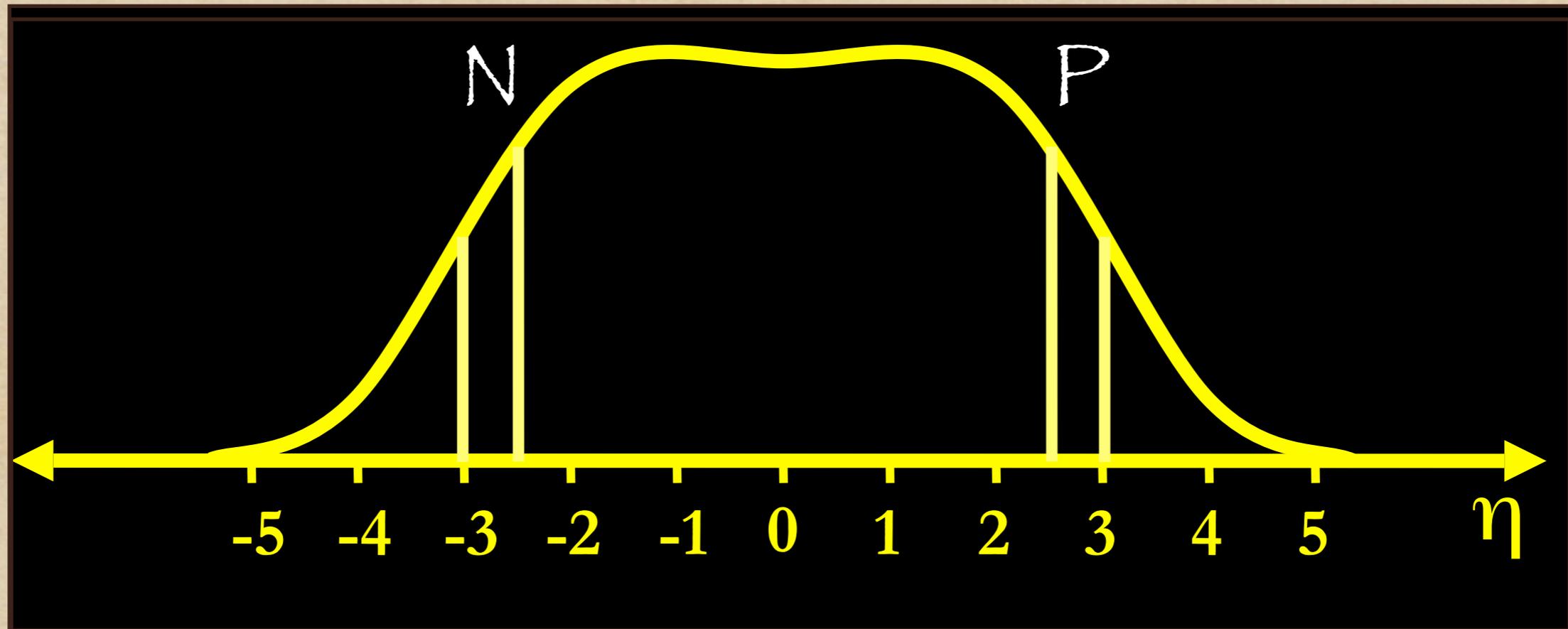
[nucl-ex/0409021](#)

# $N_{\text{part}}$ Scaling = Long Range



Integrating over  $4\pi$  gives  
simple scaling with  $N_{\text{part}}$ :  
Emphasizes non-trivial correlation between mean  
values of forward and  
mid-rapidities

# “Correlation Structure”

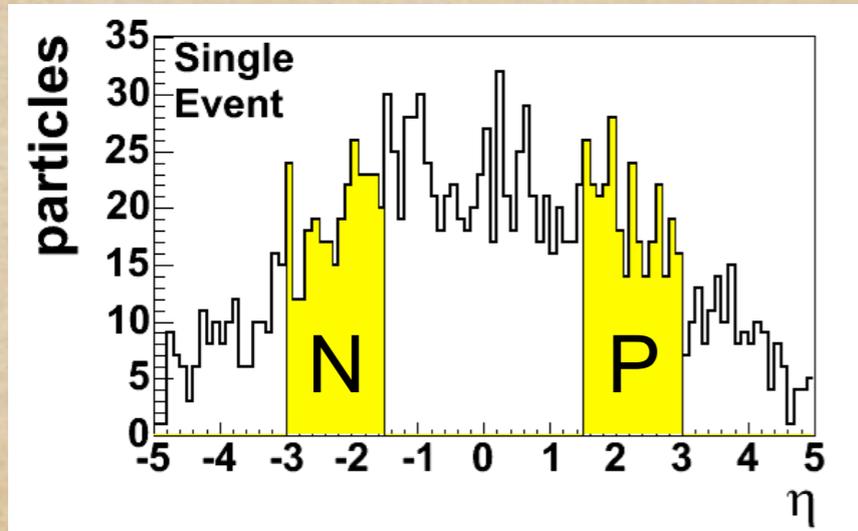


What underlies the single-particle distributions?

Are correlation short or long range?

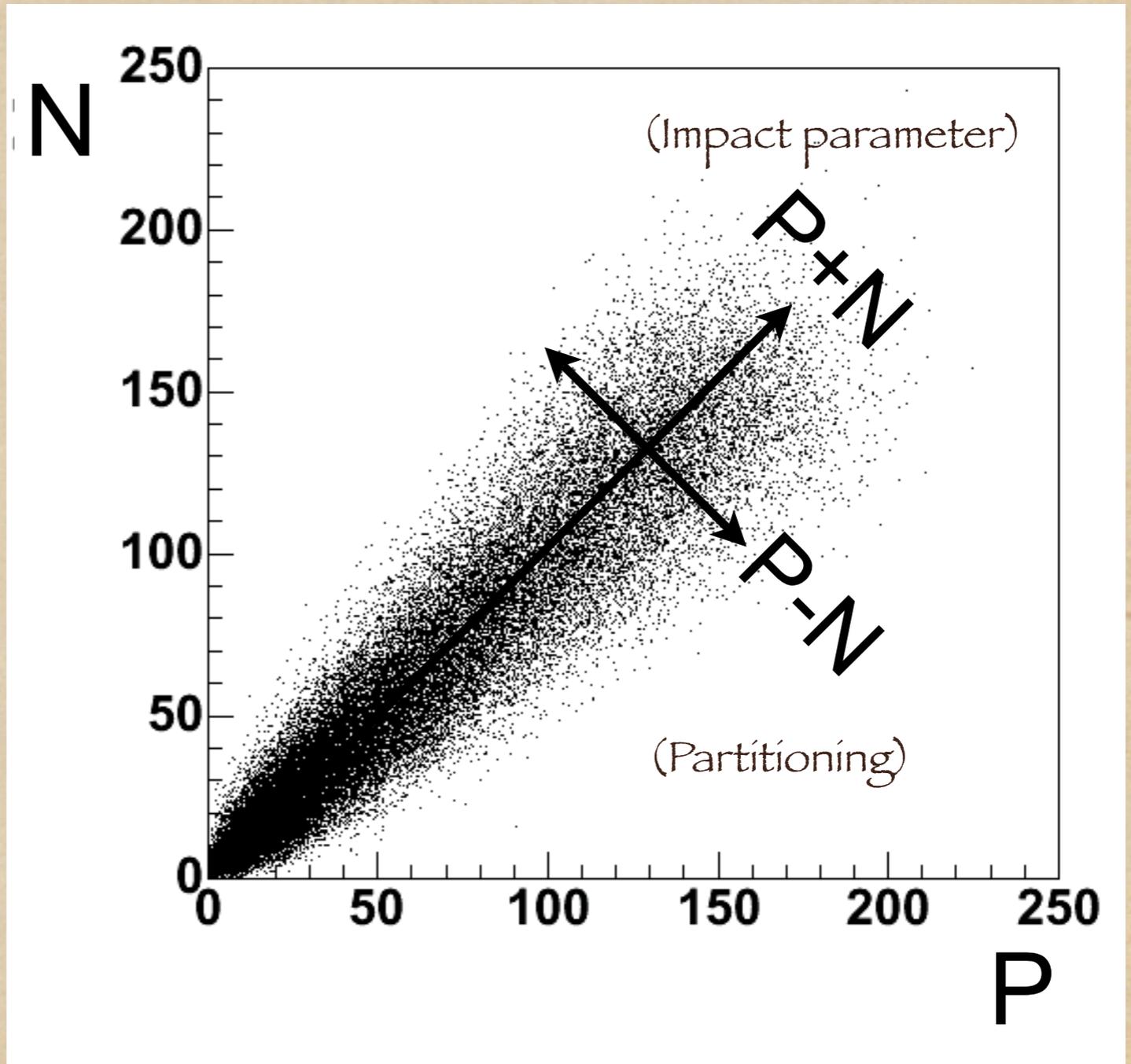
Forward-Backward Multiplicity Correlations

# Analysis Concept



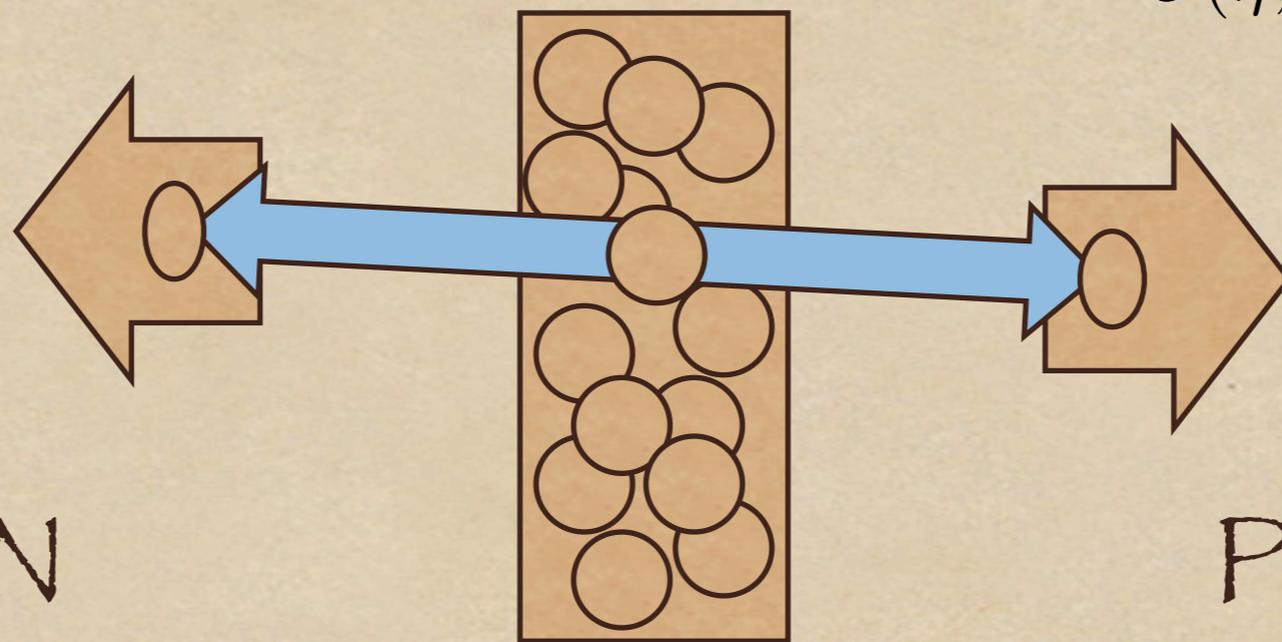
$$C = \frac{P - N}{\sqrt{P + N}}$$

RMS used to  
determine  $\sigma_C$



# Long Range Correlations

$$C(\eta, \Delta\eta) = \frac{P - N}{\sqrt{P + N}}$$



Correlated partitioning of  $X$  objects

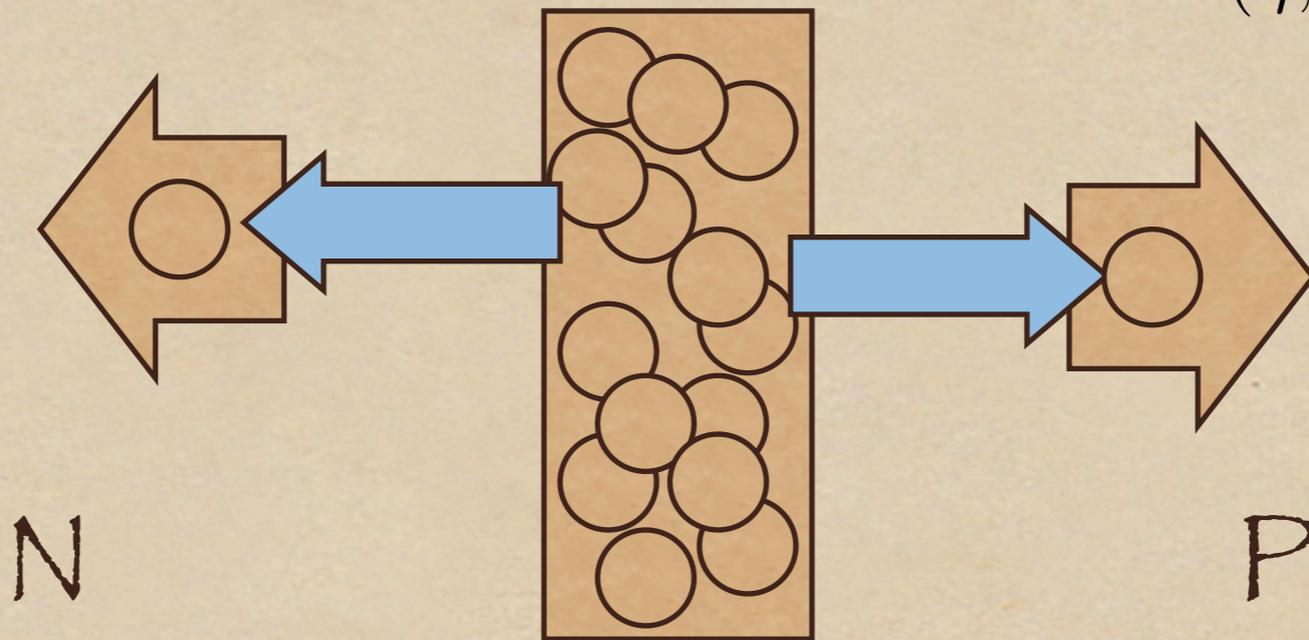
(e.g. one object splits into two)

is an “intrinsic long range correlation”

(reduces  $\sigma^2_C \rightarrow 0$ )

# “Long Range” Correlations

$$C(\eta, \Delta\eta) = \frac{P - N}{\sqrt{P + N}}$$

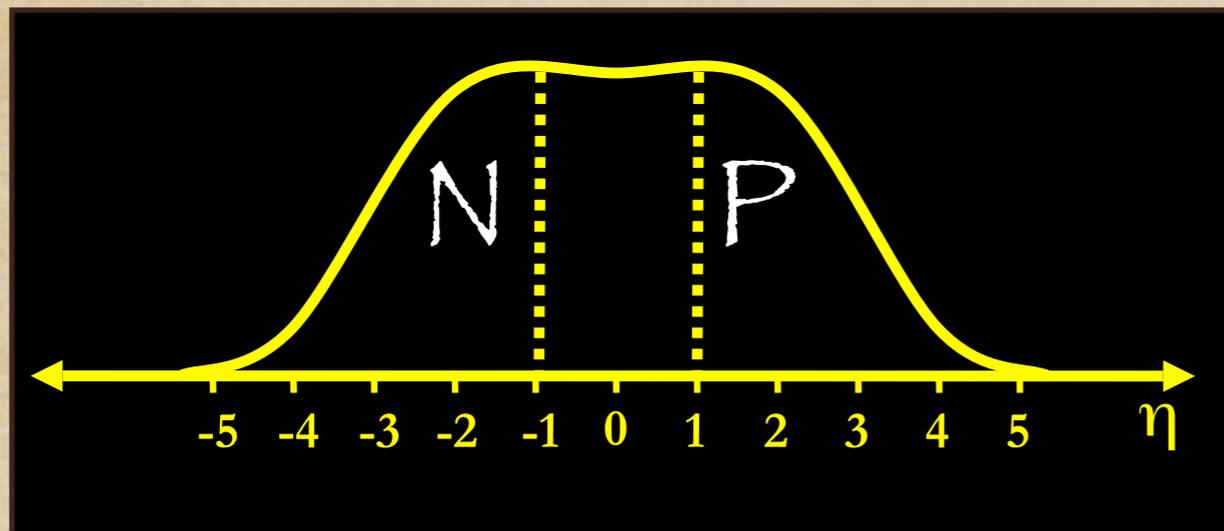


Binomial partitioning of  $X$  objects into  $P$  and  $N$  sides induces another type of “long range correlation”

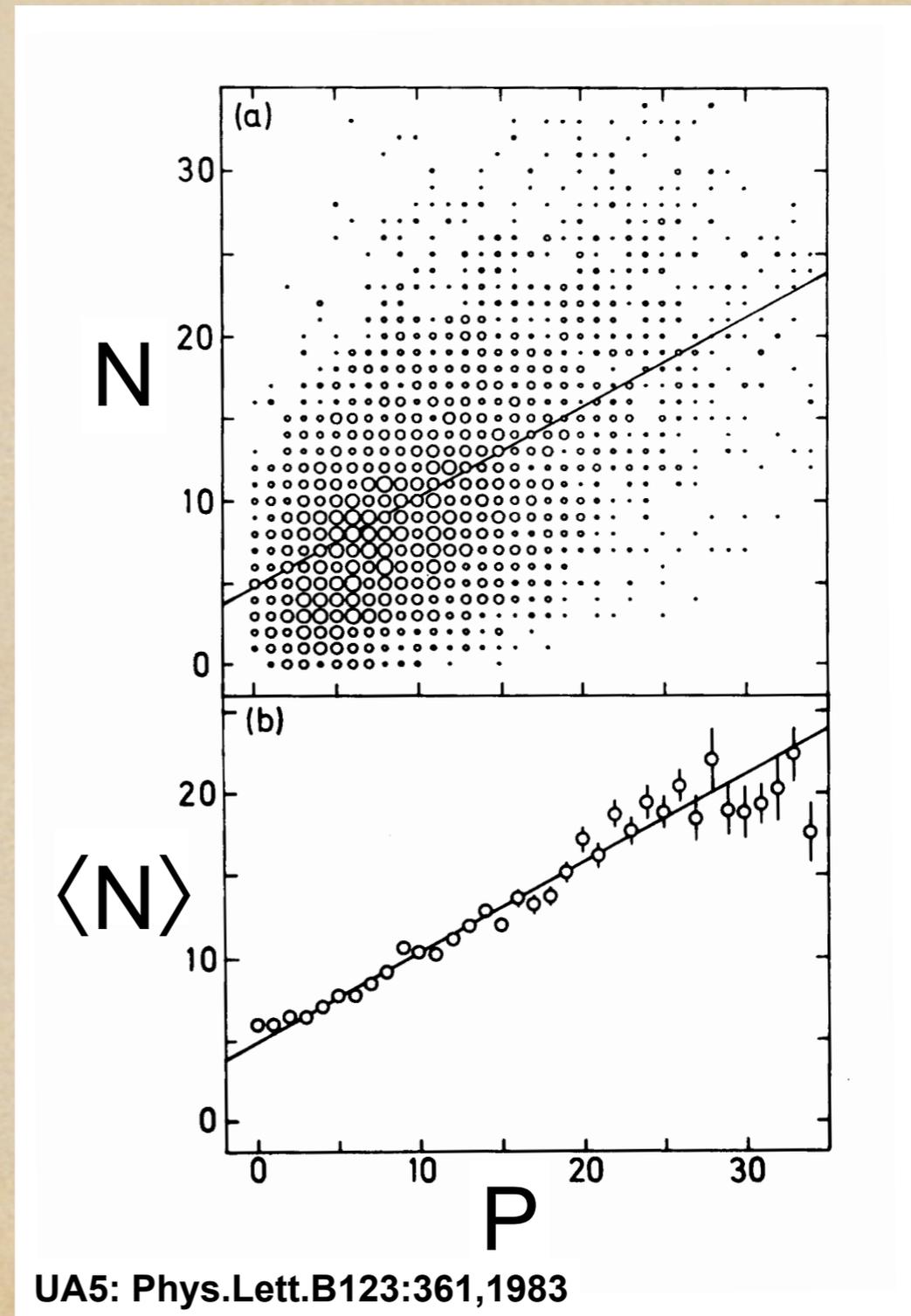
$$\sigma^2(P-N) = P+N \rightarrow \sigma^2_C = 1$$

# “Long Range” Correlations in p+p

p+p collisions have always had a “long range” component, in the binomial, i.e. non-intrinsic, sense



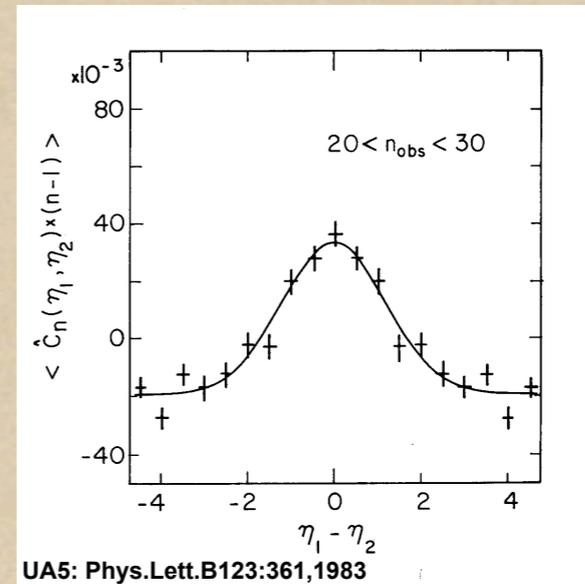
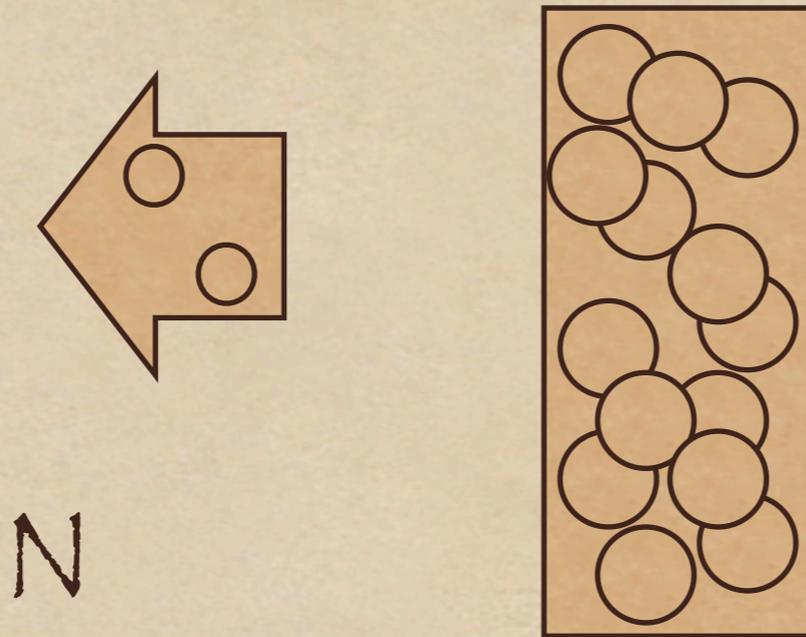
Persists when removing the central 2 units of rapidity



UA5: Phys.Lett.B123:361,1983

# Short Range Correlations

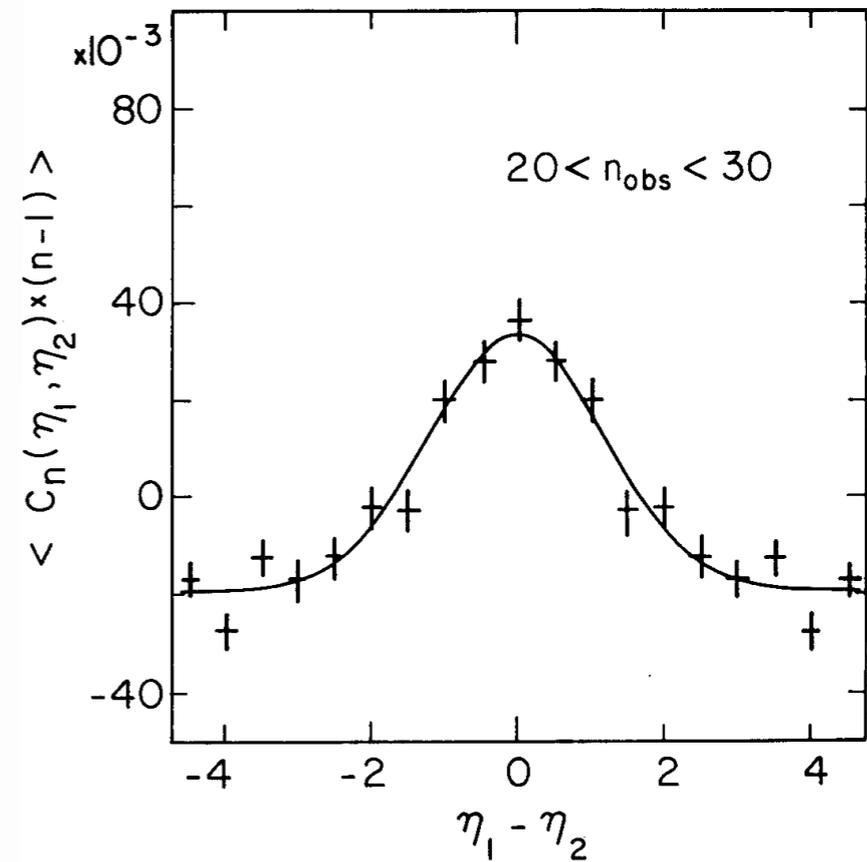
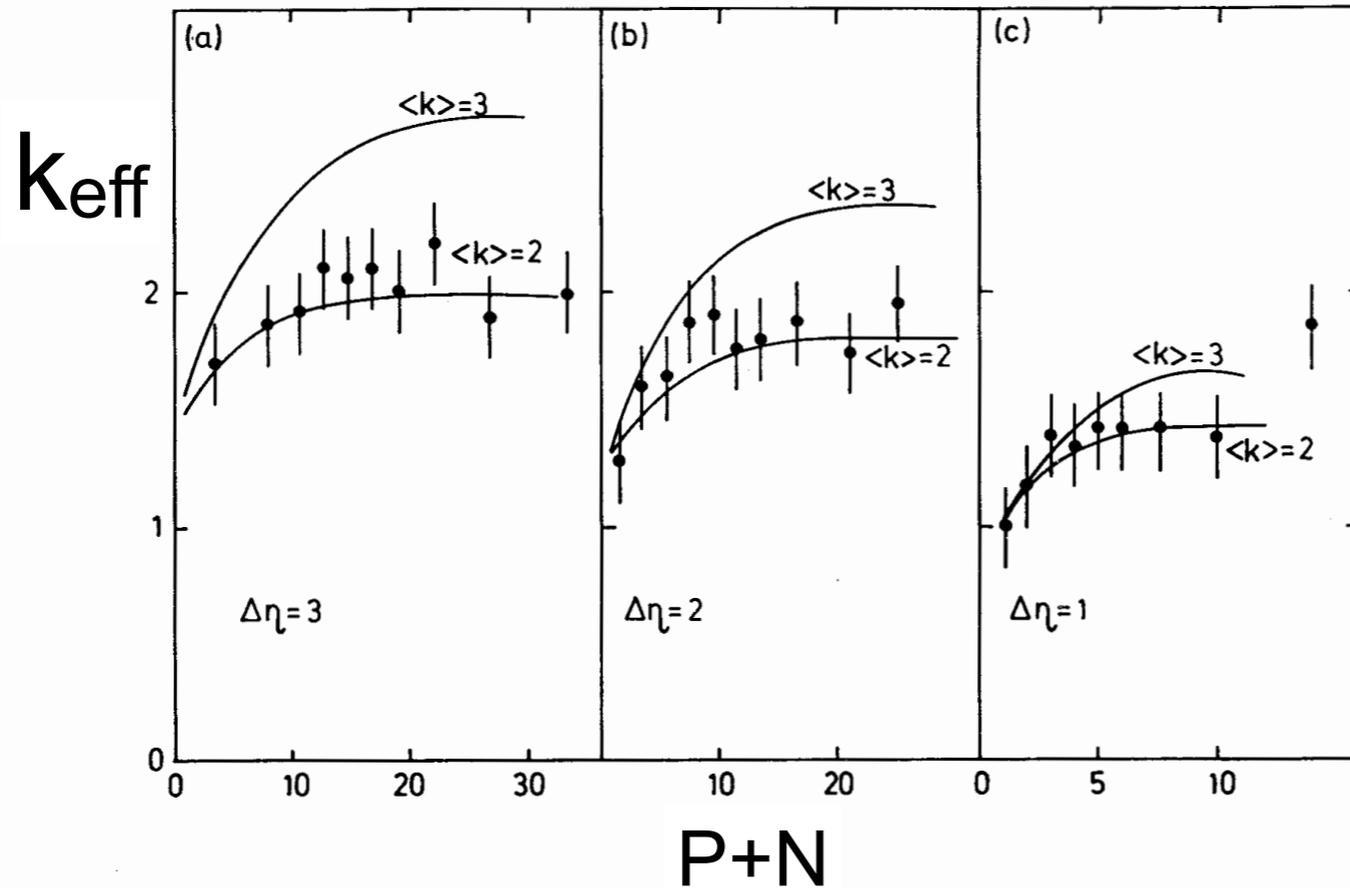
$$C(\eta, \Delta\eta) = \frac{P - N}{\sqrt{P + N}}$$



If each object breaks into  $K$  “pieces” that stay close in rapidity, induces correlations over a short range (e.g. resonances [thermal models?], gluon splitting)

# “Clusters” in $\bar{p}+p$

UA5: Phys.Lett.B123:361,1983

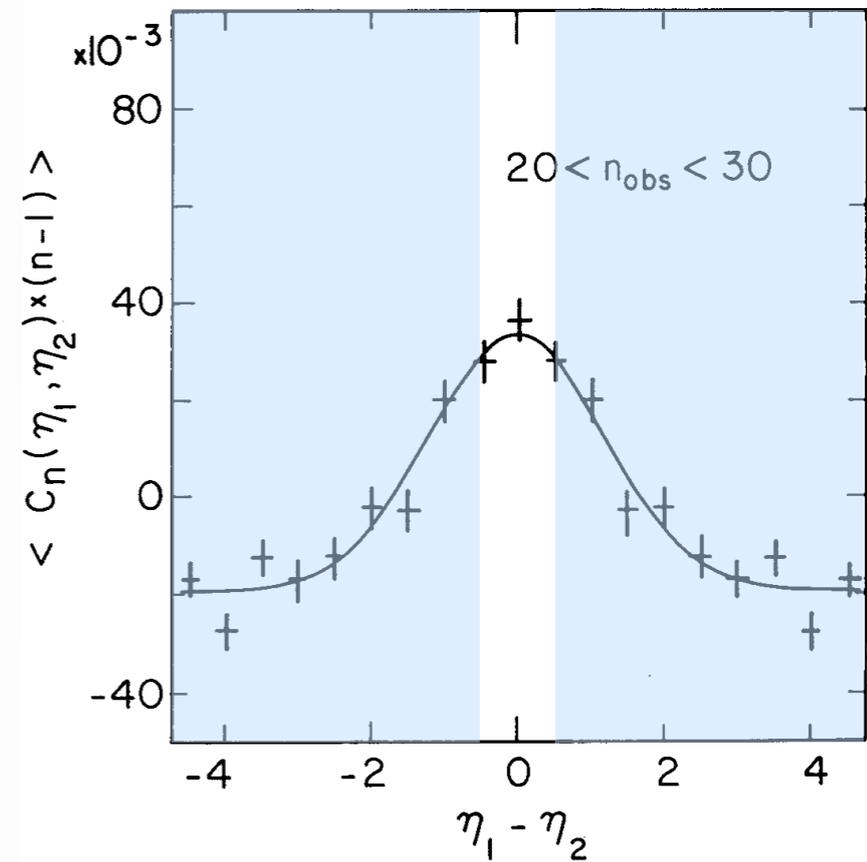
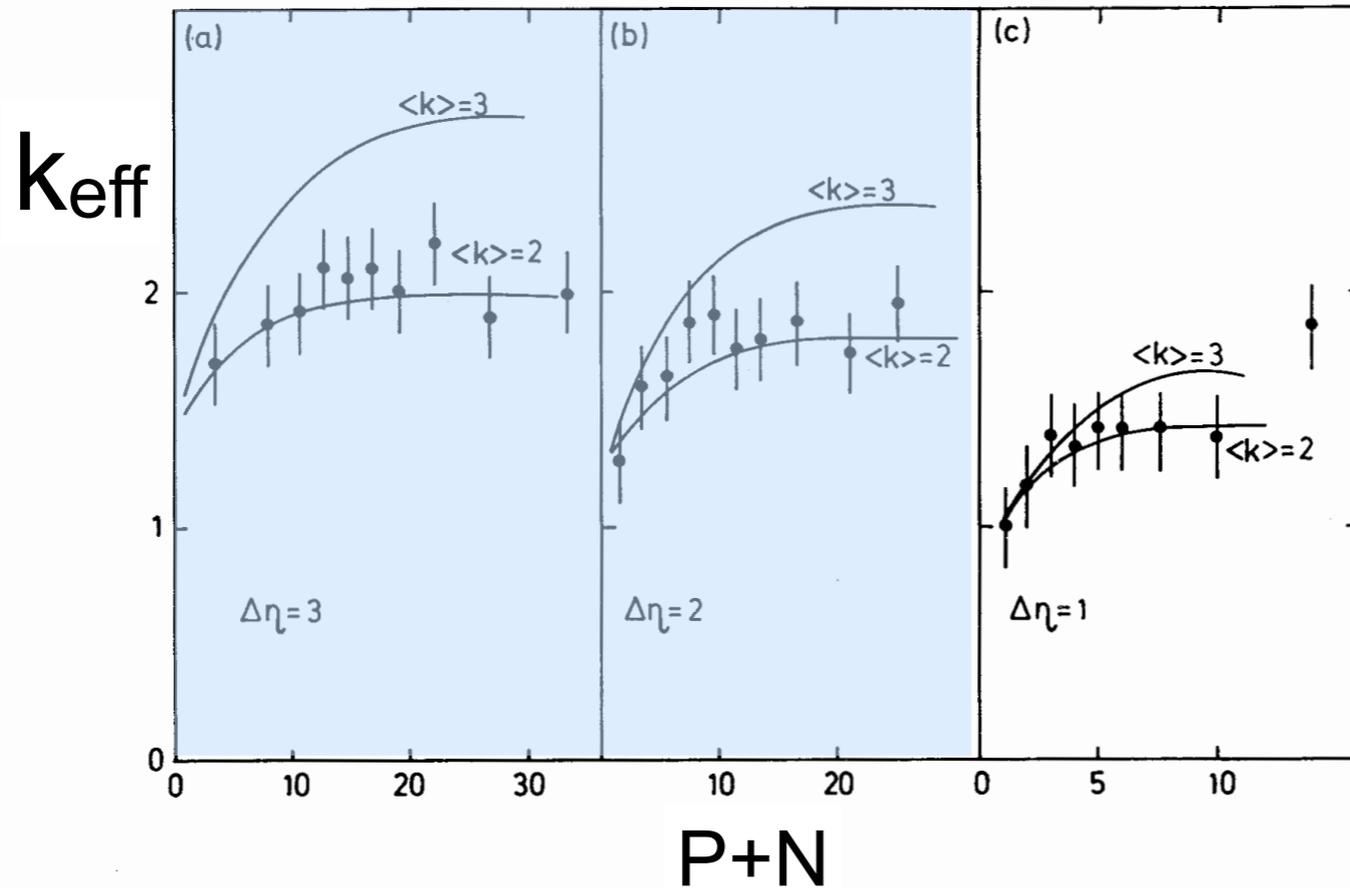


$\bar{p} + p$  546 GeV

FB correlations interpreted in terms of production of “clusters”  
 Limited rapidity windows prevent seeing all  $K$  particles  
 in a cluster  $\rightarrow$  “effective cluster multiplicity”:  $k_{\text{eff}} = \langle k \rangle + d_k^2 / \langle k \rangle$

# Clusters in $\bar{p}+p$

UA5: Phys.Lett.B123:361,1983

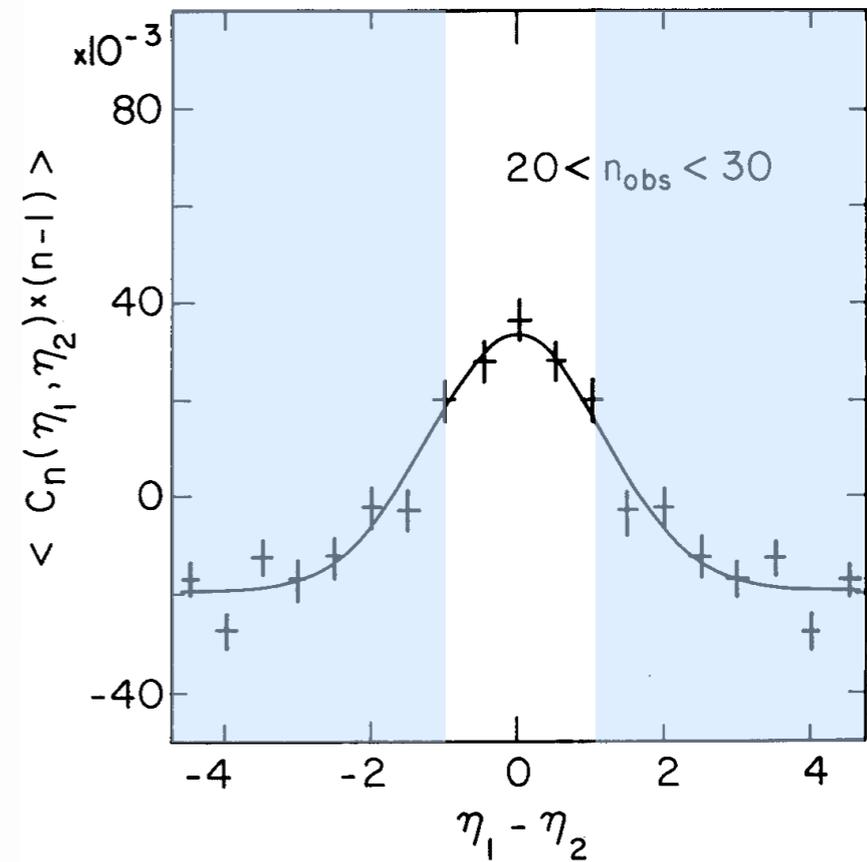
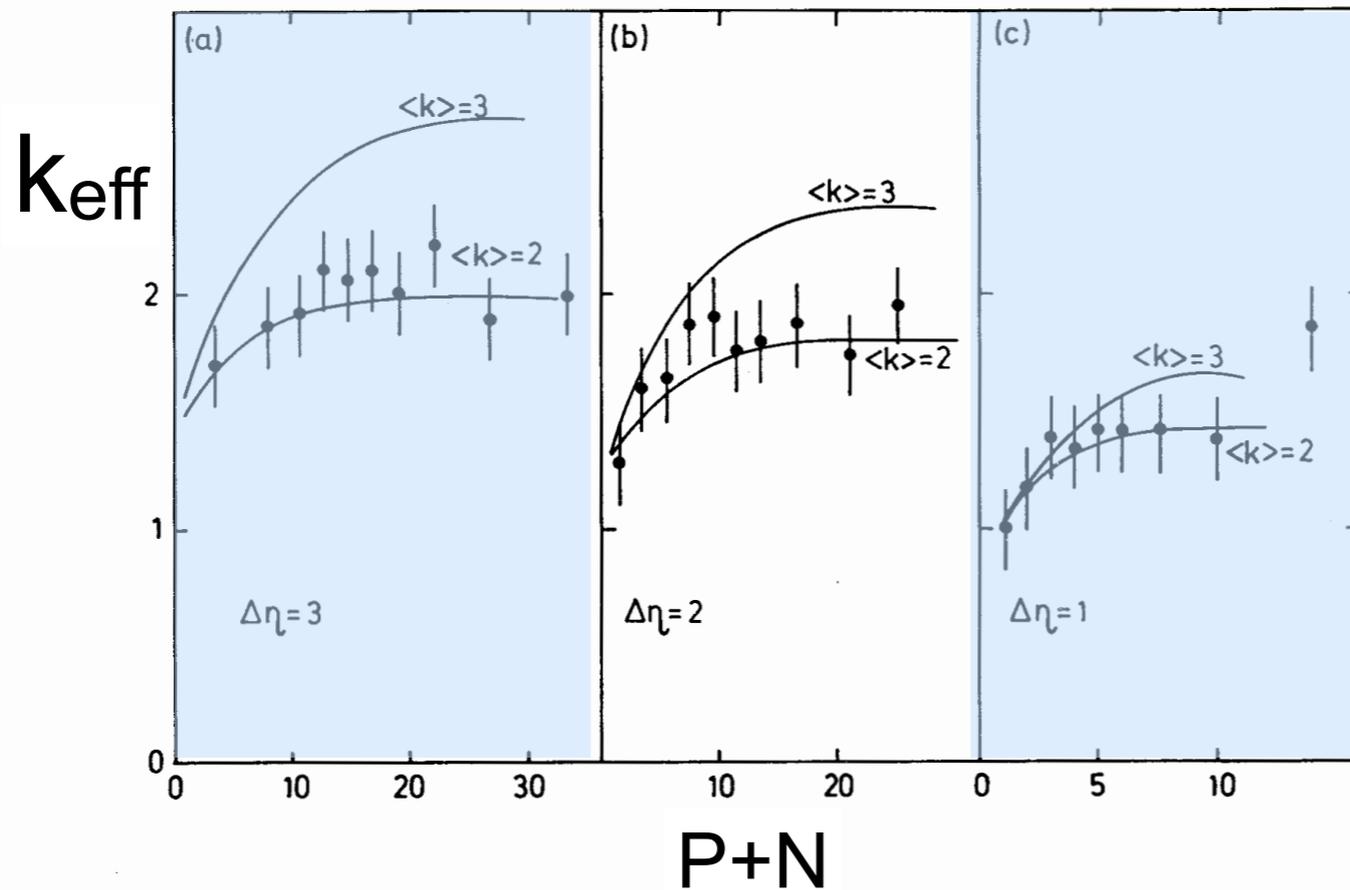


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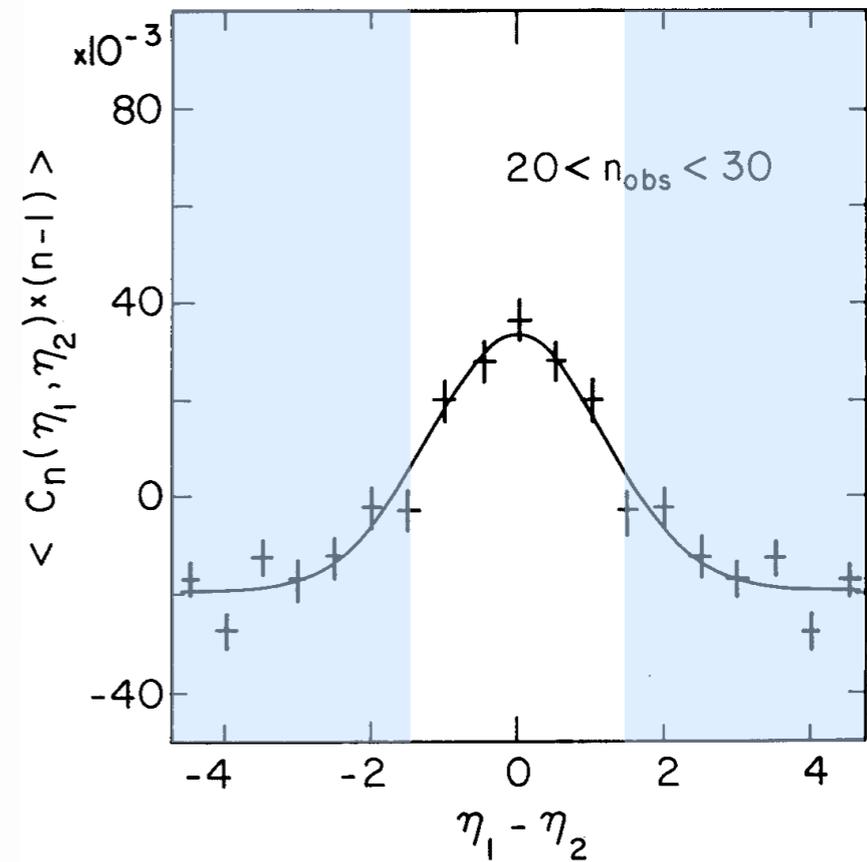
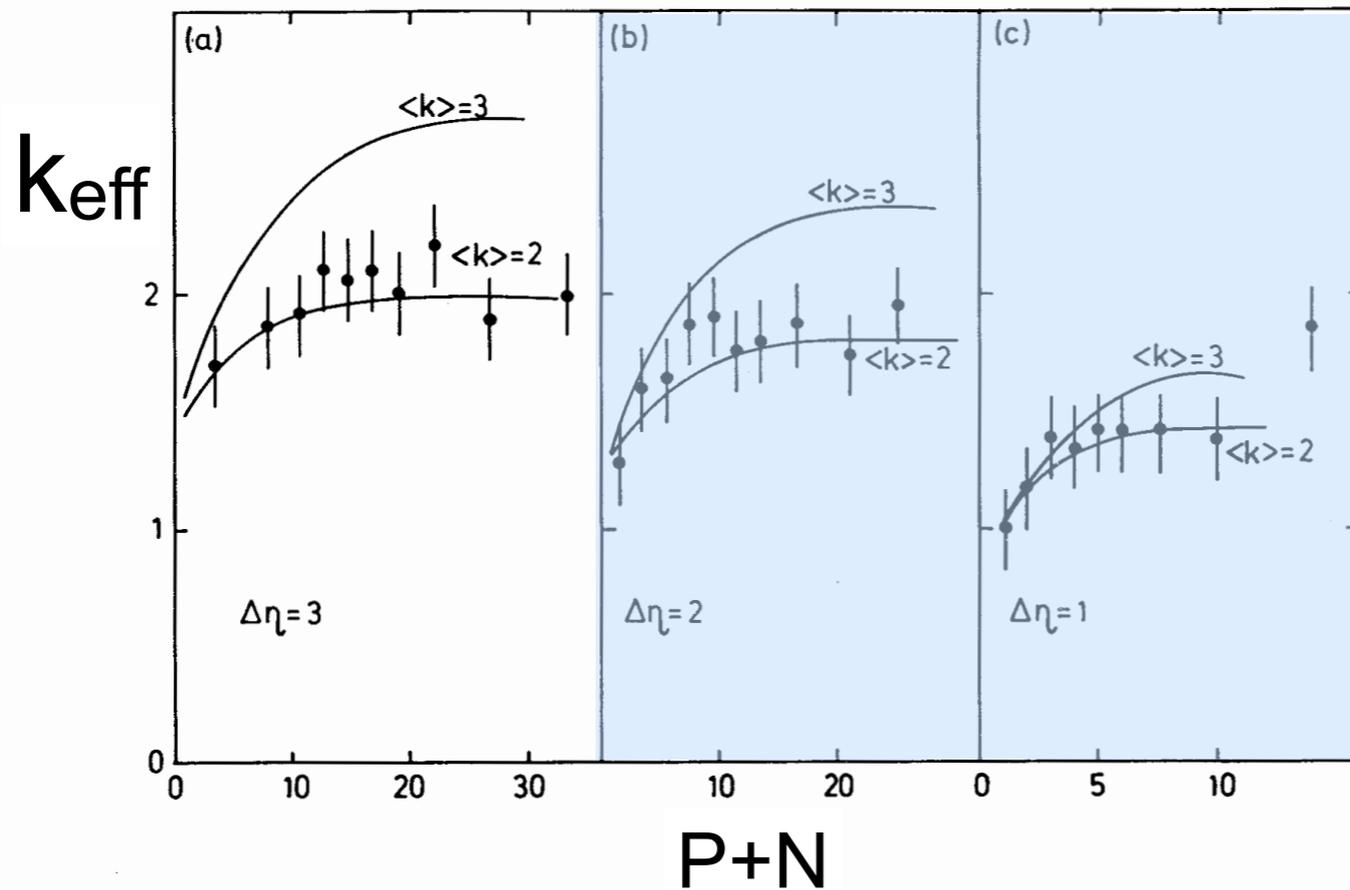


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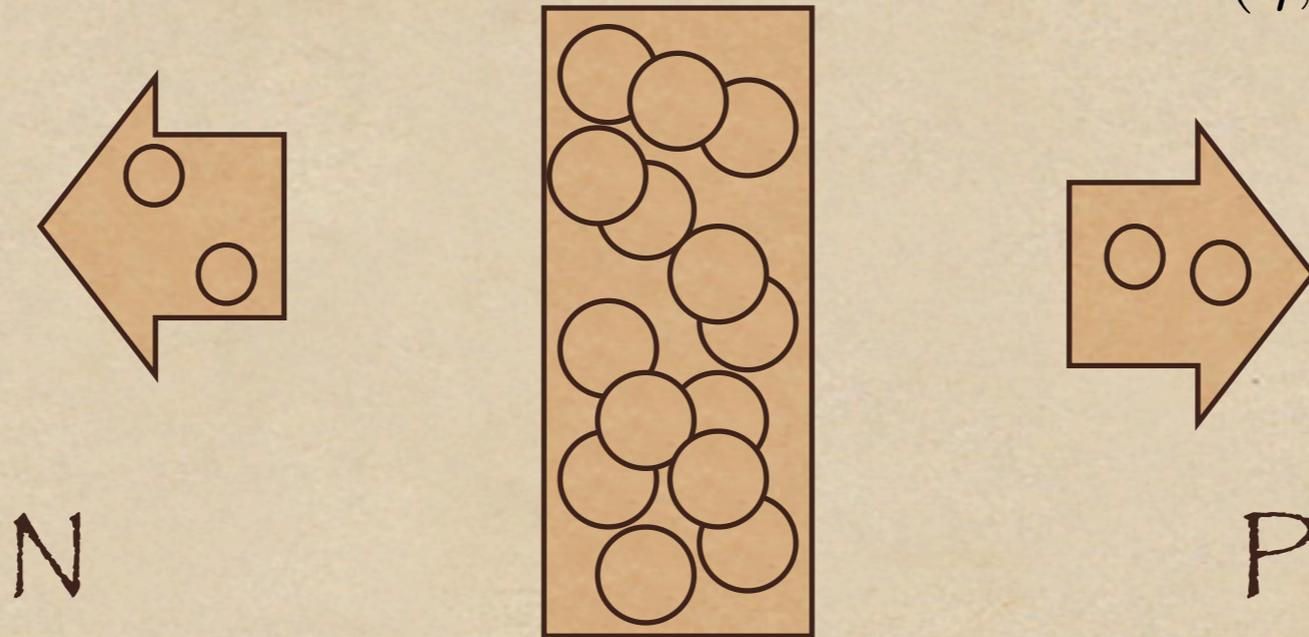


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# Clusters and $\sigma^2(C)$

$$C(\eta, \Delta\eta) = \frac{P - N}{\sqrt{P + N}}$$



$$\begin{array}{ccc}
 P \rightarrow KP & \xrightarrow{C(\eta, \Delta\eta) = \frac{P - N}{\sqrt{P + N}}} & C \rightarrow \sqrt{K}C \\
 N \rightarrow KN & & \sigma_C^2 \rightarrow K\sigma_C^2
 \end{array}$$

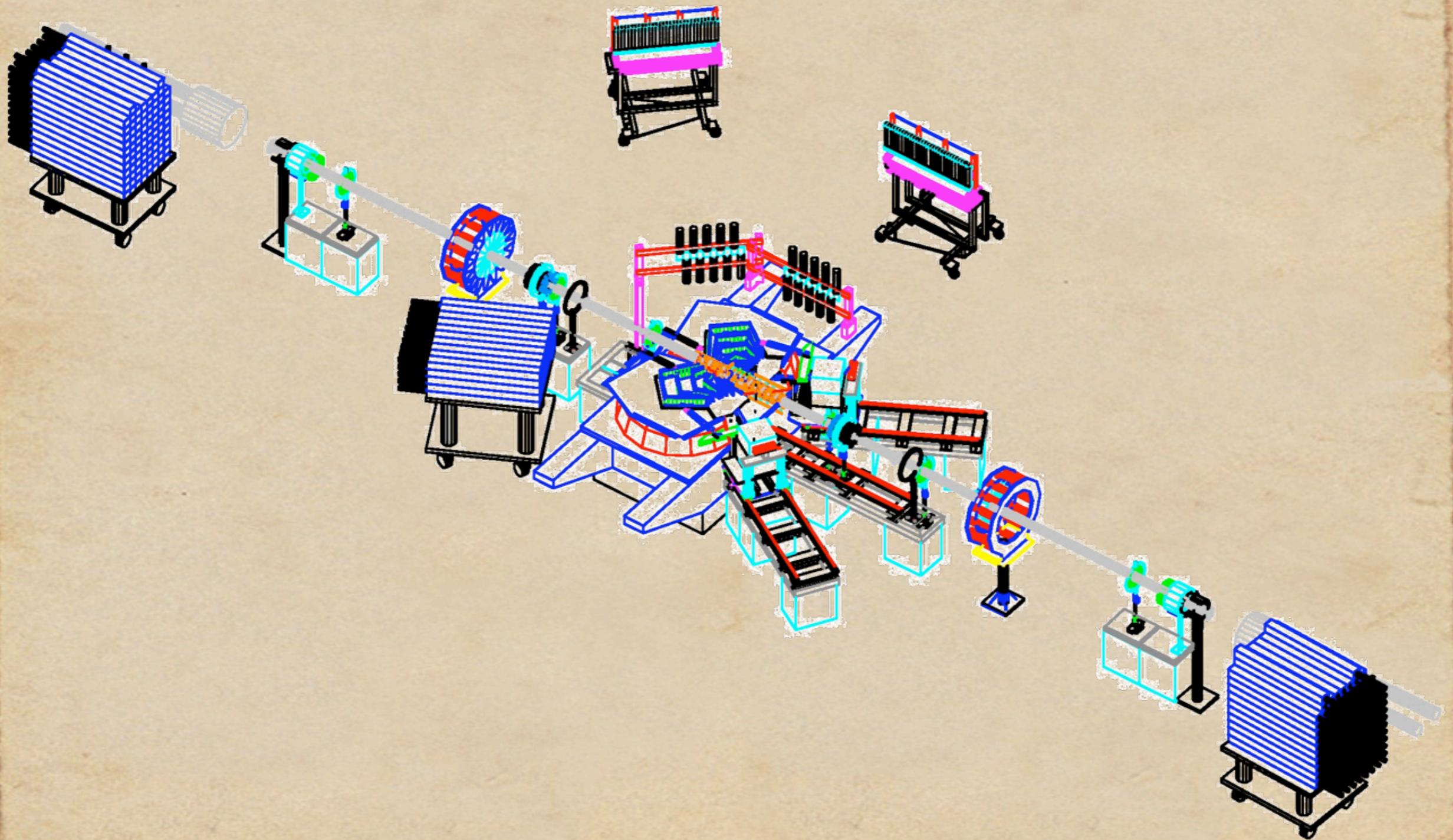
Forward Backward Multiplicity Correlations give access to cluster structure of particle production

# Summary

Intrinsic Long-Range	$\sigma^2_C \rightarrow 0$
Binomial Partitioning	$\sigma^2_C = 1$
Cluster Emission	$\sigma^2_C \propto k_{\text{eff}}$

Could expect combinations of effects

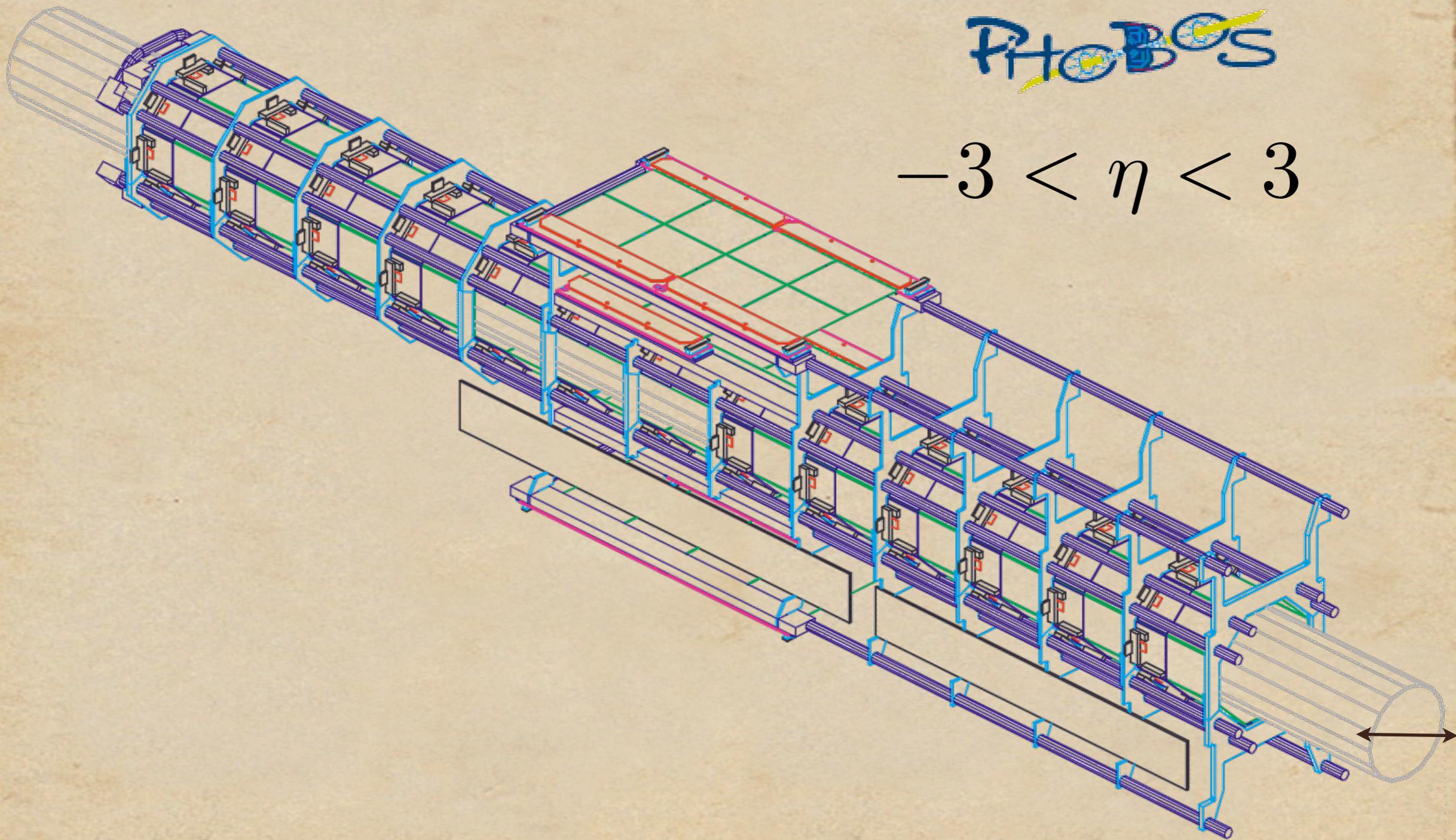
# PHOBOS 2004/5



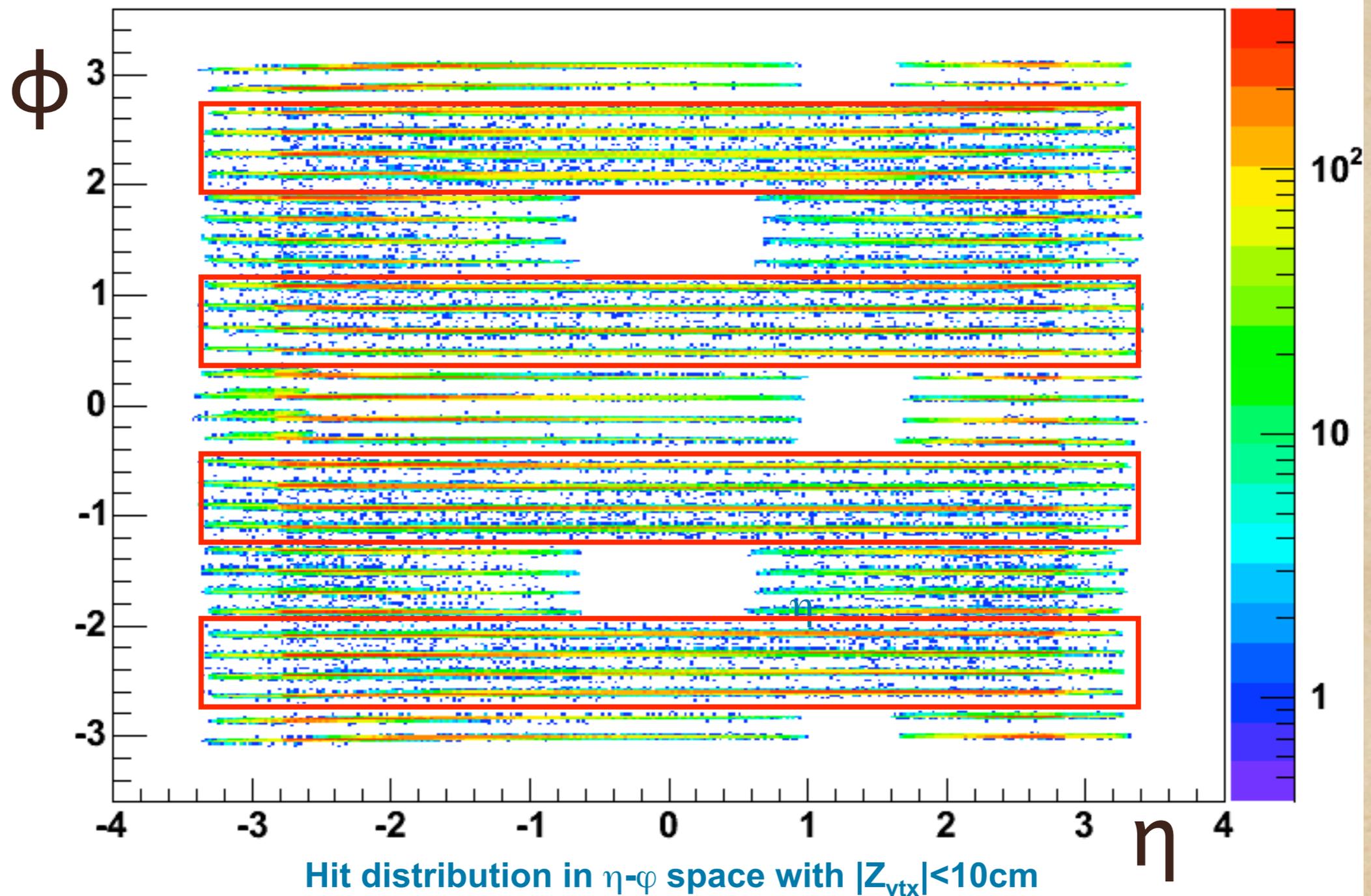
# Octagon Detector

PHOBOS

$$-3 < \eta < 3$$



# Octagon Acceptance



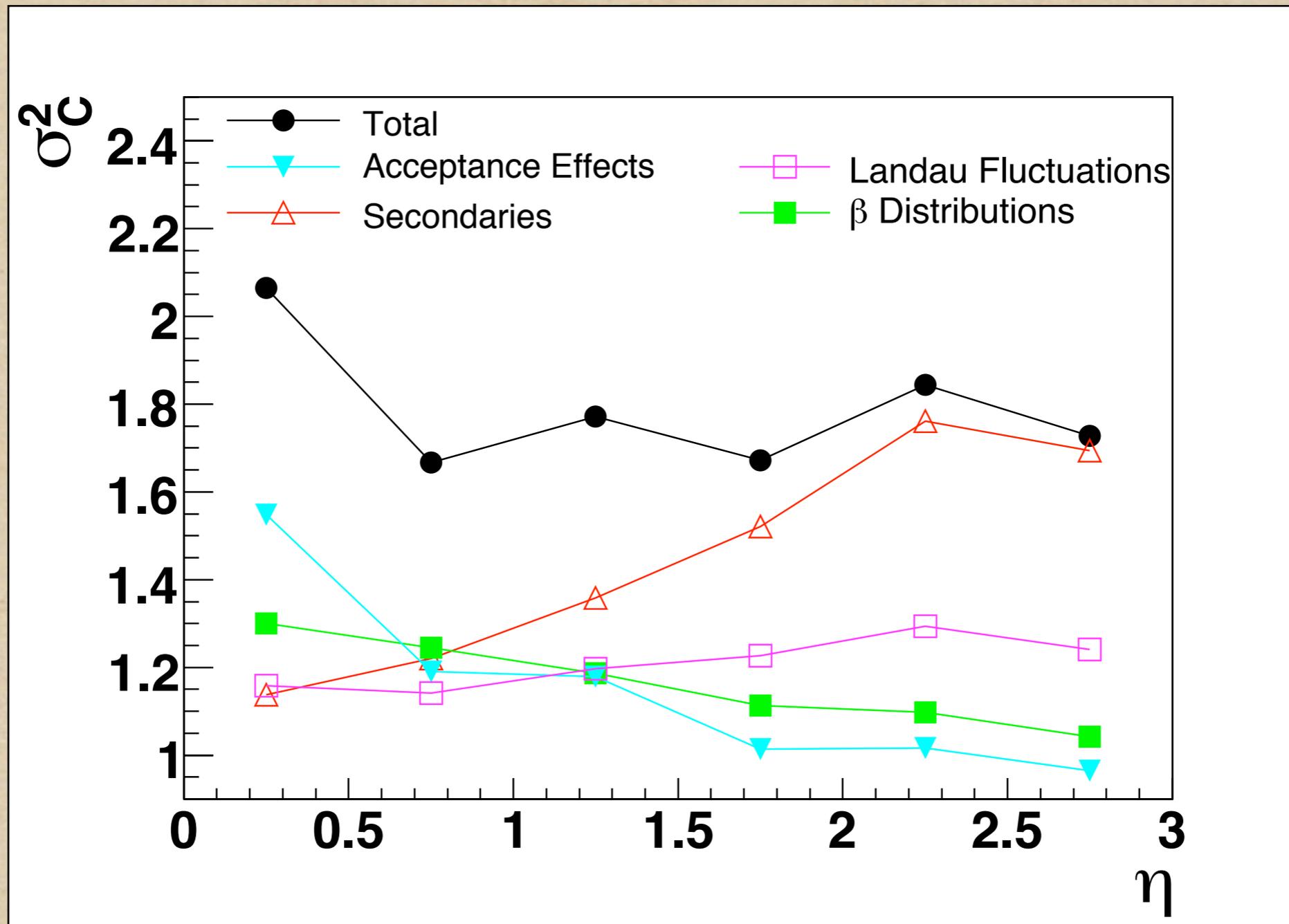
# Extraction of $\sigma^2_C$

$$\sigma^2_{C,raw} = \sigma^2_C + \cancel{\sigma^2_{det}}$$
$$\rightarrow k_{eff}$$

Removing detector effects gives access to this effective cluster size (modified by acceptance / bin-width)

NEW for QM2005!

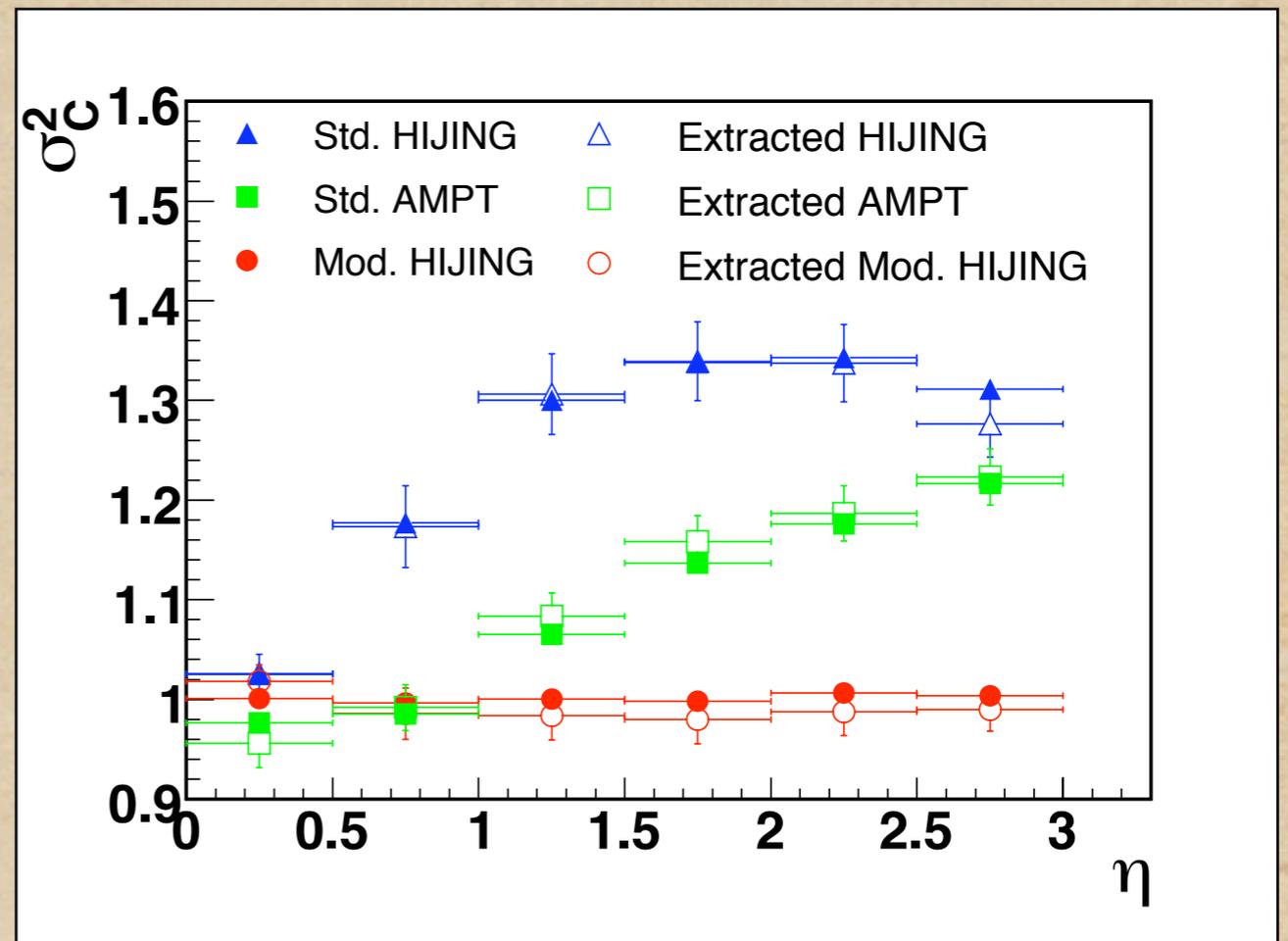
# Detector Effects vs. $\eta$



Different sources contribute flat in  $\eta$

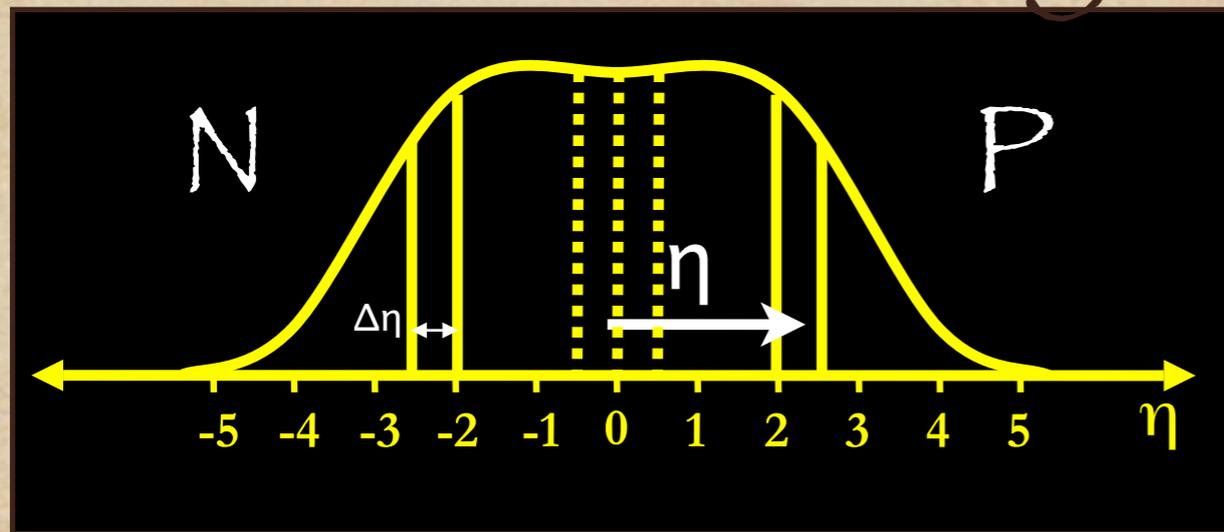
# Corrections & Systematics

- ◆ Residual detector effects:  
 $\sigma^2_{\text{det}}$
- ◆ Acceptance gap effects
- ◆ Secondaries contribution at large  $\eta$
- ◆ Half  $\rightarrow$  Full azimuth

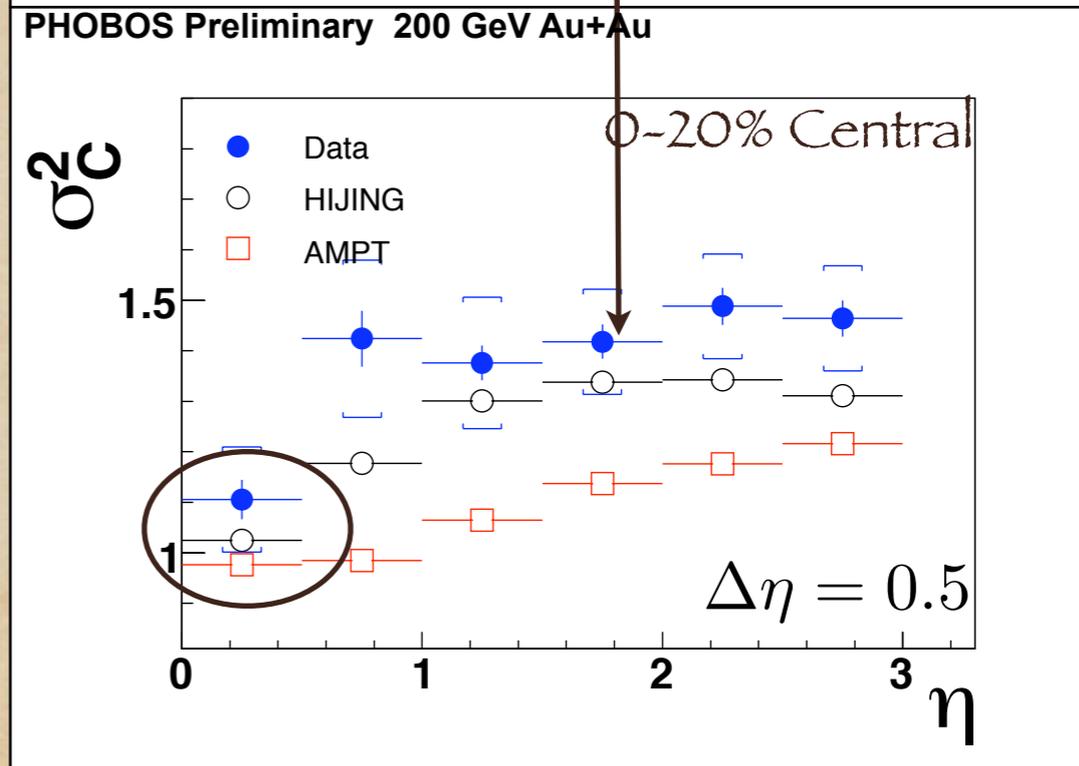
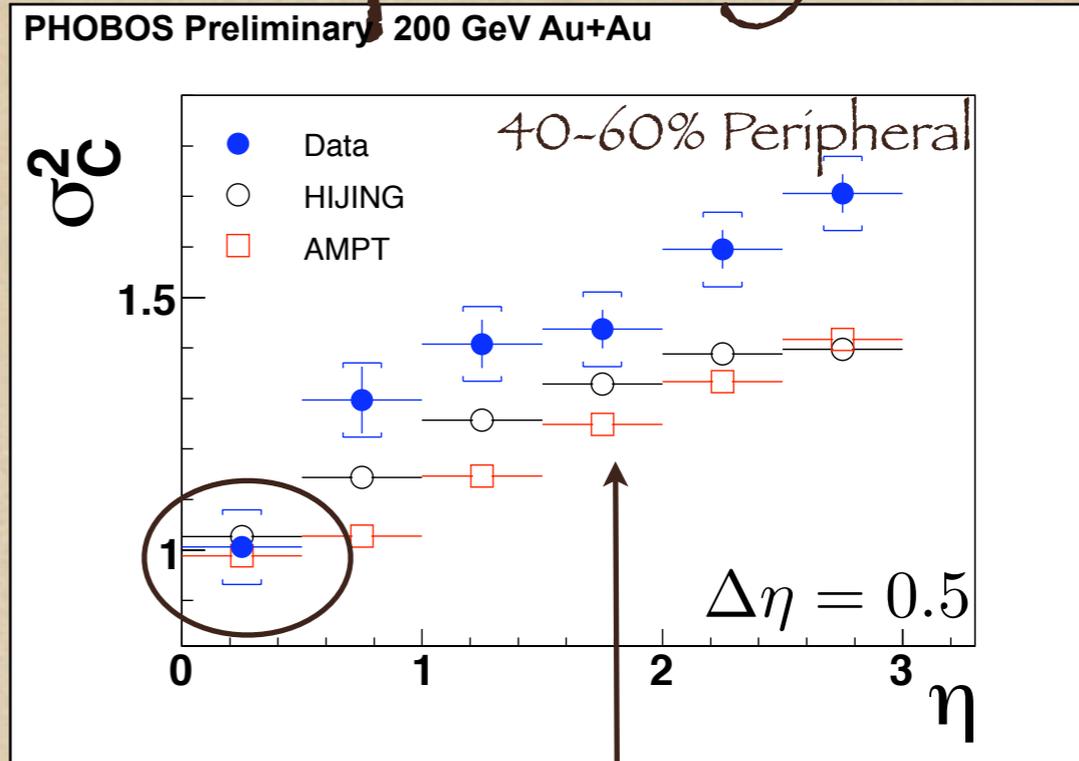


Systematics calculated for several sources:  
use average value of combined error  $\Delta\sigma^2_C \sim 0.1$

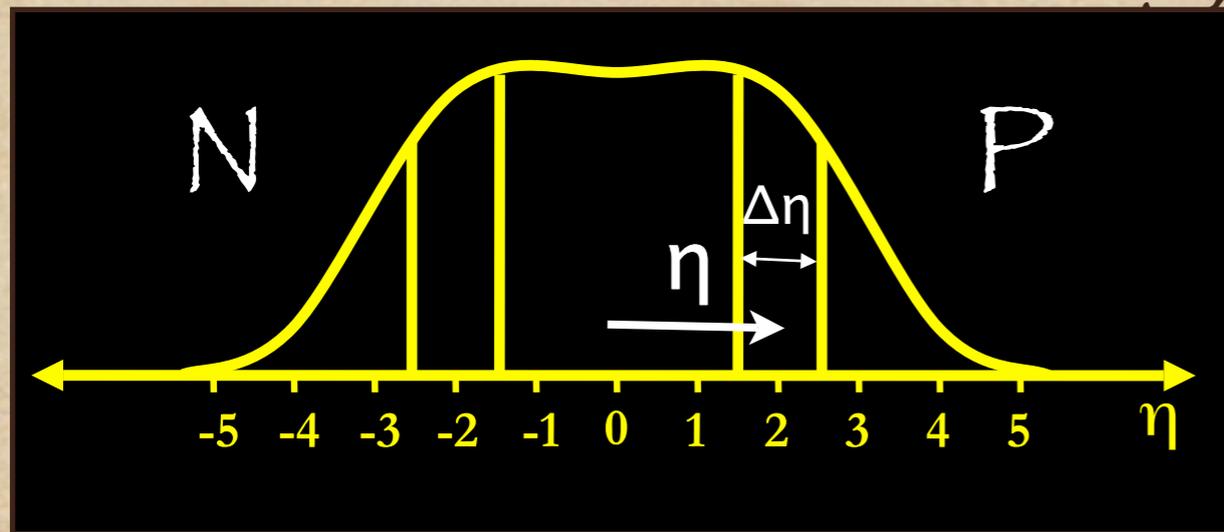
# Centrality & Rapidity



- \* HIJING & AMPT agree in peripheral events, diverge in central events
- \* Significant centrality dependence
- \* Suppression at  $\eta=0$  from intrinsic correlation of P&N



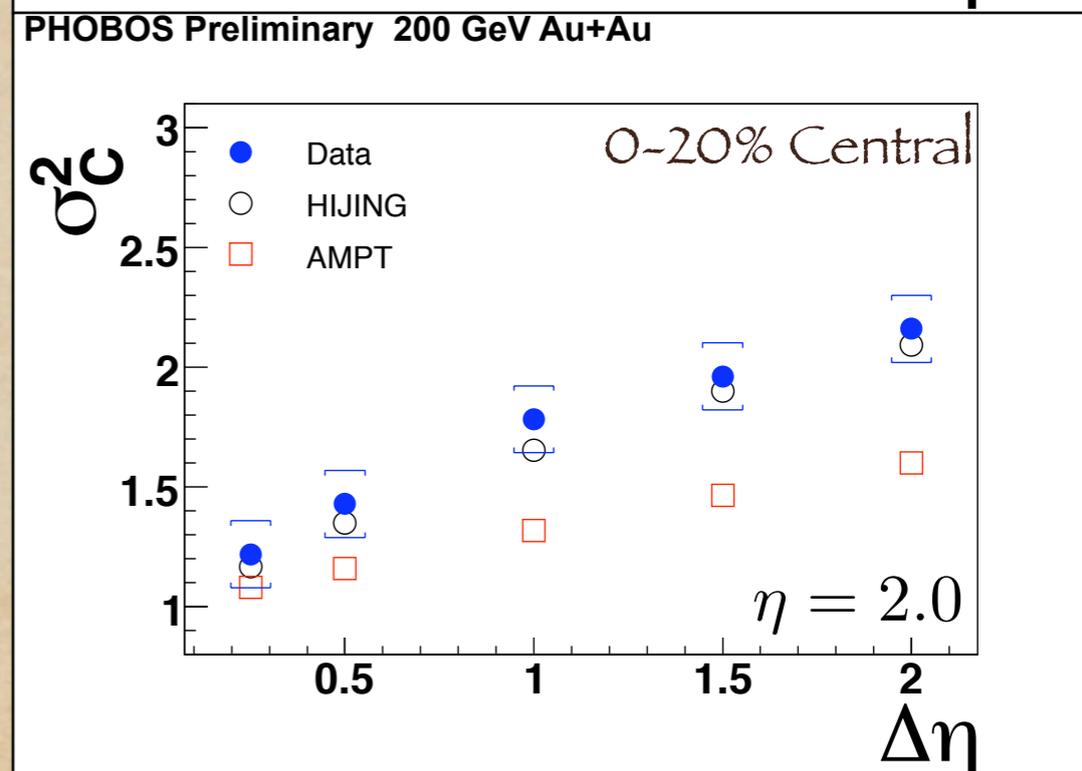
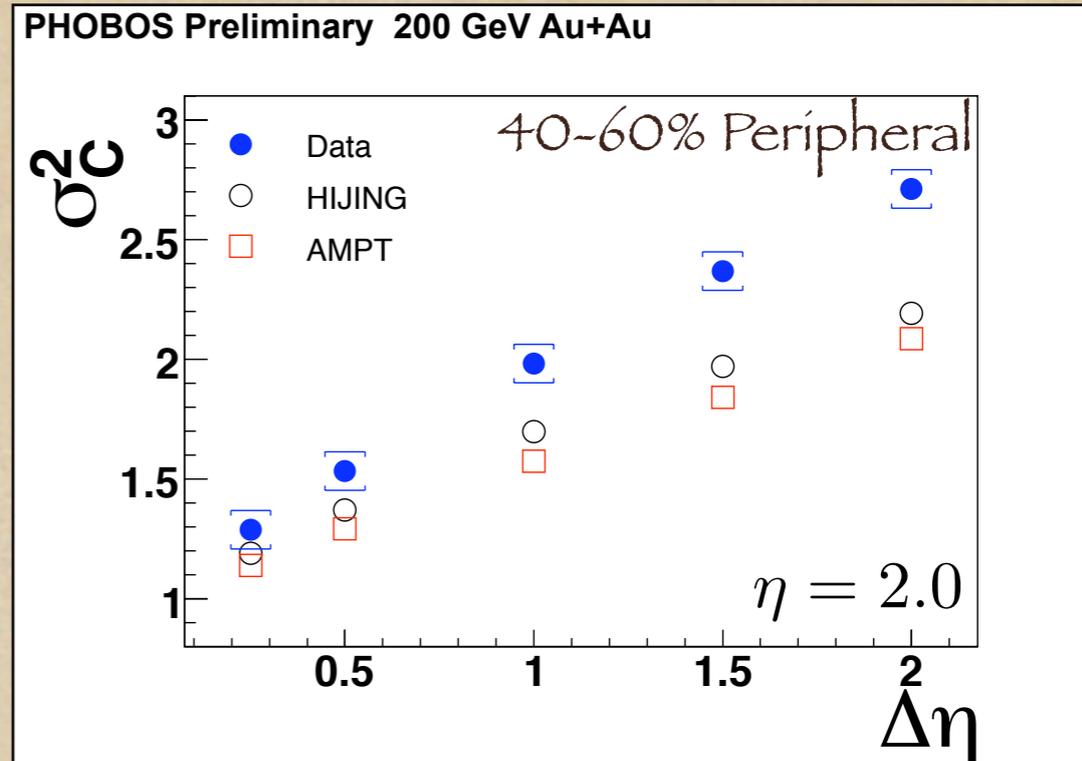
# Centrality & Width



\* Rate of change with bin size, reflects full cluster multiplicity  $K$

\* Peripheral data implies larger  $K$  than HIJING & AMPT

\* HIJING ~ Central data;  
AMPT ~ Centrality dependence



# Summary

- ◆ Forward-Backward multiplicity fluctuations provide information about long and short range effects in pseudorapidity
- ◆ New PHOBOS data for 200 GeV Au+Au with detector effects corrected
  - ◆ Clear short-range correlations observed
  - ◆ Non-trivial centrality and rapidity dependence
  - ◆ Central data similar to HIJING; centrality dependence qualitatively similar to AMPT

# Charge Fluctuations & QGP

Hadron gas creates  
multi-particle clusters  
(defined as what  
is seen in p+p)

QGP should “smooth”  
out fluctuations,  
decreasing effective  
cluster size near  $y=0$

Studies by Jeon, et al

nucl-th/0503085

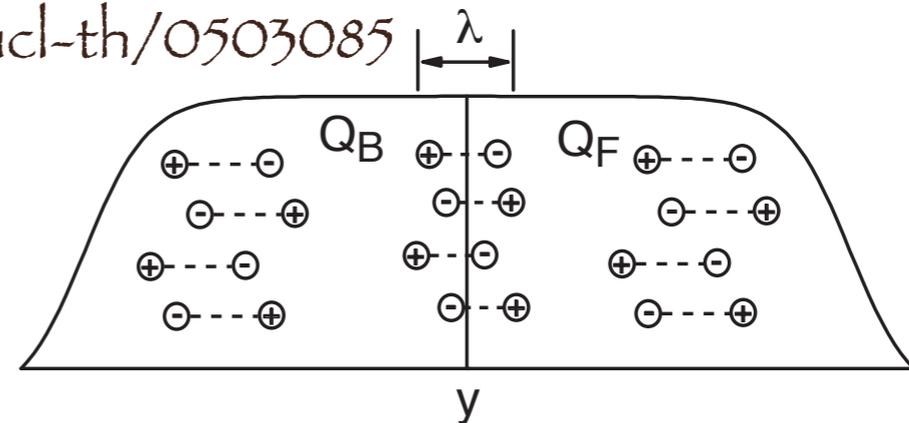


FIG. 1: A schematic illustration of the charge transfer fluctuations in the rapidity space. Only the pairs within  $\lambda/2$  of  $y$  can contribute to the charge transfer fluctuation  $D_u(y)$ . Here  $\lambda$  is the rapidity correlation length, or the rapidity distance of the decay particles from a single cluster. If  $\lambda$  is a function of  $y$ , then  $D_u(y)$  also changes with  $y$ .

PHOBOS data should  
provide limits on such  
a scenario



# Systematic Errors vs. $\eta$

