Equation of state from lattice QCD

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Outline

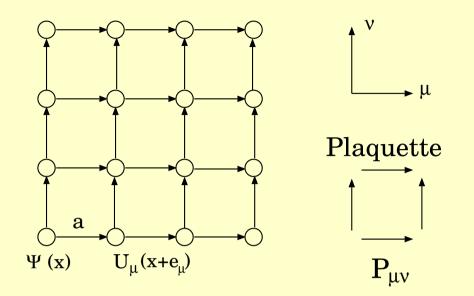
- 1. Introduction
- 2. Recent results
- 3. New results for the Equation of State
- 4. Summary

Introduction

• The Equation of State (EoS); p, ϵ, s as a function of T is an unambiguous prediction of the QCD Lagrangian

- The EoS is an important input for hydrodynamical models of heavy-ion collisions
- Perturbation theory is only reliable at very large T
- Lattice QCD is an applicable non-perturbative tool to determine the EoS

Lattice QCD introduction



Fundamental Fields:

Gauge fields:

 $U_{\mu}(x) \in SU(3)$ live on the links

Quarks:

 $\Psi(x), \bar{\Psi}(x)$

anti-commuting Grassmann variables live on the sites

Wilson fermions: O(a) artefacts Staggered fermions: $O(a^2)$, BUT flavour symmetry violation Partition function

$$Z = \int dU d\Psi d\bar{\Psi} e^{-S_E}$$

 S_E is the Euclidean action

Parameters:

gauge coupling g

quark masses m_i ($i = 1..N_f$)

(Chemical potentials μ_i)

Volume (V) and temperature (T)

Finite $T \leftrightarrow$ finite temporal lattice extension

$$T = \frac{1}{N_t a}$$

Continuum limit: $a \rightarrow 0$

Renormalization: keep the physical spectrum constant

at finite T:

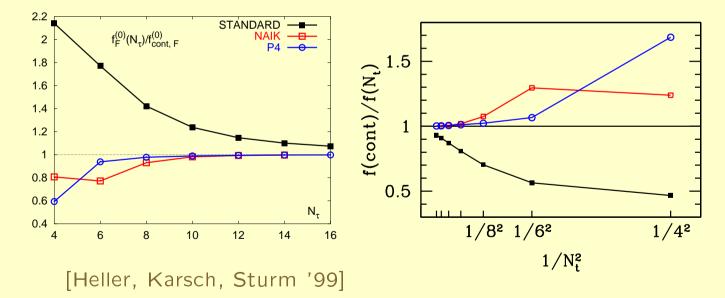
continuum limit $\iff N_t \to \infty$

Improved actions

 S_E is not unique; many possibilities

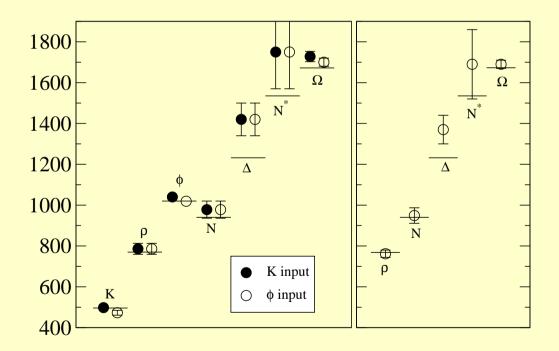
from flic to clover and tadpole, hyper-improved and even overimproved improvements

Continuum limit is always important!



Continuum extrapolation from N_t and $N_t + 2$ standard action may be better than using only N_t with improved action How reliable is lattice QCD?

At T = 0: Hadron spectrum based on the QCD Lagrangian (quarks+gluons) no more – no less than the experimental spectrum quantitative agreement on the percent level already in the quenched approximation



[Hasenfratz, Juge, Niedermayer, 2004]

At T > 0: no clear connection between experiments and lattice (yet) experiences from T = 0 are promising

Equation of state from lattice simulations

energy density (ϵ) and pressure (p) from partition function:

$$\epsilon(T) = \frac{T^2}{V} \frac{\partial(\log Z)}{\partial T} \qquad \qquad p(T) = T \frac{\partial(\log Z)}{\partial V}$$

T, V are varied by a, take derivative with respect of a

$$\frac{\epsilon - 3p}{T^4} = -\frac{L_t^3}{L_s^3} a \frac{d(\log Z)}{da}$$

the pressure $(p \propto \log[Z])$ along the LCP by the integral method:

$$\frac{p}{T^4} = L_t^4 \int d(\beta, m \cdot a) \left(\frac{\partial (\log Z)}{\partial \beta}, \frac{\partial (\log Z)}{\partial (m \cdot a)} \right)$$

Renormalization of the pressure

We want p(T = 0) = 0 and $\epsilon(T = 0) = 0 \rightarrow$ Simulations at both

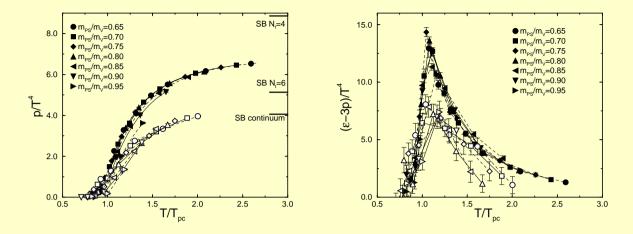
T > 0 ($N_t \ll N_s$) and T = 0 ($N_t \gtrsim N_s$) are necessary and then subtraction:

$$\frac{p}{T^4} = \frac{p_T}{T^4} - \frac{p_0}{T^4}; \qquad \frac{\epsilon}{T^4} = \frac{\epsilon_T}{T^4} - \frac{\epsilon_T}{T^4}$$

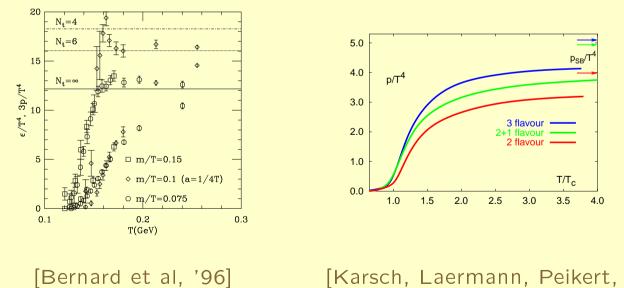
numerical precision needed for the subtraction increases with $N_t^4 \rightarrow \text{CPU}$ costs grow faster $(\mathcal{O}(1/a^{13}))$ than for T = 0 simulations

Today $N_t = 4$ is easy $N_t = 6$ is difficult $N_t = 8$ is a challenge

Recent lattice results Wilson fermions: $\mathcal{O}(a)$, slower



[Ali-Khan et al, '01] Staggered fermions: $\mathcal{O}(a^2)$, faster



[Karsch, Laermann, Peikert, 2000]

Ongoing projects: MILC, Bielefeld-Brookhaven-Columbia

- 1. Unrealistic quark masses might be important, since $T_c \ge m_{\pi}$
- 2. No Line of constant physics (LCP) used $T = 1/(N_t a)$ is increased with decreasing aphysical spectrum $(m_{\pi}, m_K, m_{\rho}, ...)$ should not change
- 3. flavour symmetry violation (staggered) unphysical, large pion non-degeneracy
- Approximate algorithms were used
 R algorithm: systematic error due to finite stepsize
 high precision subtraction can be sensitive to it
- 5. Lattice artefacts improved action with $N_t = 4$ only
- 6 Scale determination no string-tension in dynamical QCD

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 $N_t = 4, 6$

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 $q\bar{q}$ force at 0.5fm

 $N_t = 4, 6$

New lattice results for the EOS

[Y. Aoki, Z. Fodor, SDK, K.K. Szabo]

Main features:

- Physical mass spectrum is used for T > 0 simulations
- Use of LCP:

physical spectrum unchanged while *a* changes

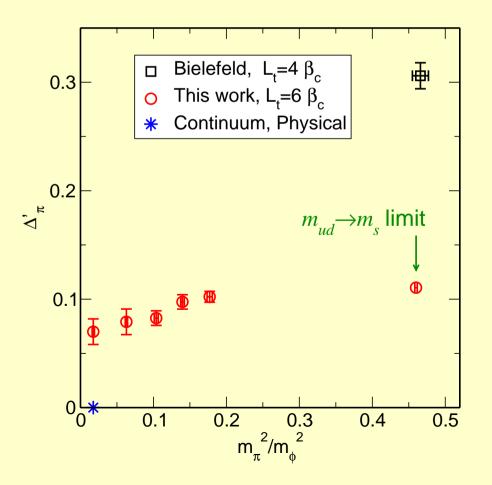
- Exact algorithm (RHMC) is used to get rid of stepsize errors
- Supressed flavour symmetry violation
 1-loop improved Symanzik gauge action +
 stout improved fermionic action
- Two sets of lattice spacings $N_t = 4$ and 6 simulations
- Unambiguous scale setting

Stout improvement

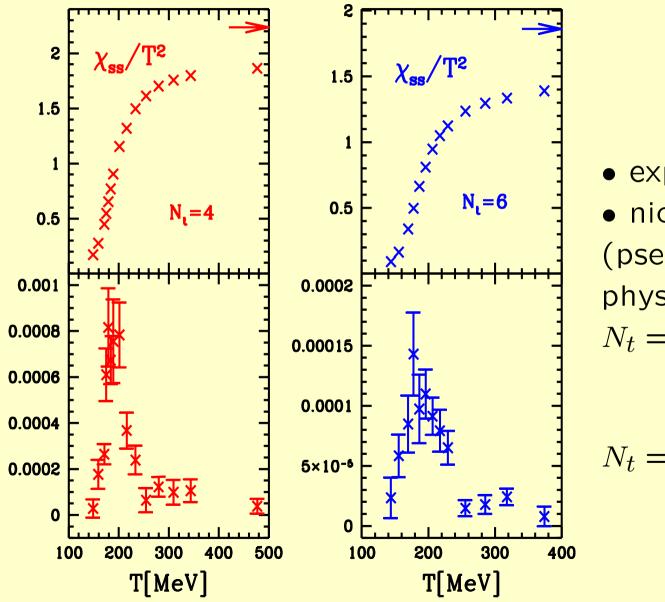
Stout smearing: replace the $U(x)_{\mu}$ gauge links with V stout links:

$$V = P \left[\rightarrow +\rho \left(\sqrt{\gamma} + \sqrt{\gamma} + \sqrt{\gamma} \right) \right]$$

unphysical non-degeneracy of pions largely reduced:



Quark number susceptibilities transition temperature



 $\chi_{ff'} = \frac{T}{V} \frac{\partial^2 \log Z}{\partial \mu_f \partial \mu_{f'}}$

• experimentally relevant

• nice peak in $\partial \chi_{ss} / \partial \beta$ (pseudo)critical coupling for physical quark masses: $N_t = 4$:

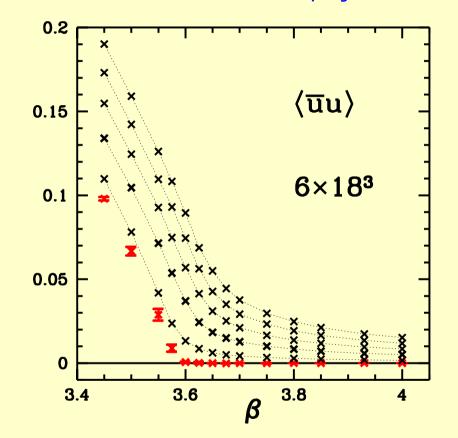
 $T_c = 186(3)(3)$ MeV

 $N_t = 6$:

 $T_c = 193(6)(3)$ MeV

Chiral condensate

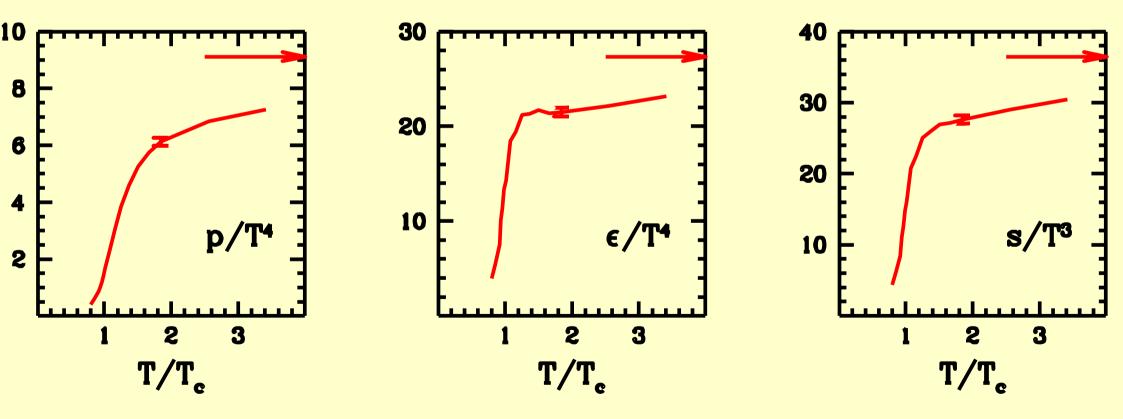
Simulations for $m_{ud} = \{1, 3, 5, 7, 9\}m_{phys}$ at finite T



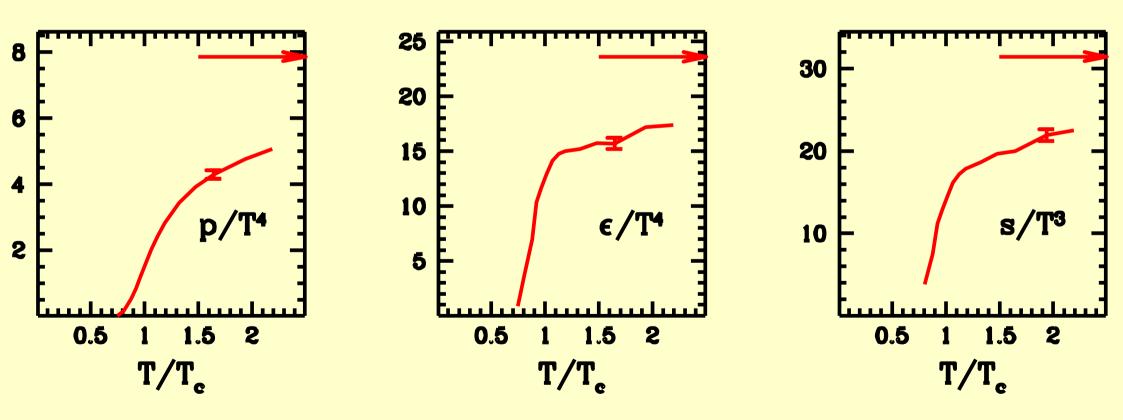
extrapolate to $m = 0 \rightarrow 2^{nd}$ order phase transition expected

 $N_t = 6$: $T_c(m = 0) = 191(5)(2)$ MeV

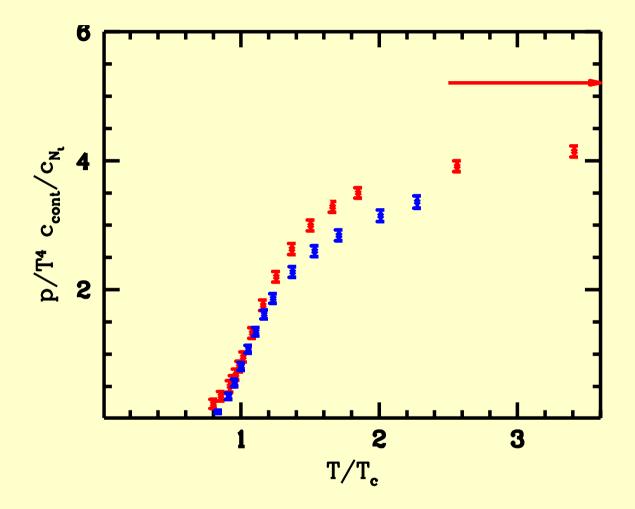
The pressure, energy density and entropy density for $N_t = 4$



The pressure, energy density and entropy density for $N_t = 6$



Scaling of the pressure Comparison of $N_t = 4$ and $N_t = 6$



• No good scaling yet. Most probably $N_t = 4$ is too coarse $\rightarrow N_t = 8$ might be needed for final continuum-extrapolated result

Summary, Conclusions

- Previous results on EoS suffer from several weaknesses
- New results improve on these points
- Transition temperature using different methods: $T_c \approx 189(8)$ MeV
- EoS is presented for two sets of lattice spacings
- Continuum-extrapolation already possible, but better to wait for even finer lattices