

Centrality dependence of heavy flavor production from single electron measurement in $\sqrt{s_{NN}}=200$ GeV Au+Au collisions

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We present preliminary measurements of electron production in p+p, d+Au, and Au+Au collisions at $\sqrt{s_{NN}}=200$ GeV for transverse momenta $1.5 \text{ GeV}/c < p_T < 8 \text{ GeV}/c$ as a function of centrality. These measurements were carried out using the STAR Time Projection Chamber and Barrel Electromagnetic Calorimeter. In this manuscript we describe the measurement techniques used to discriminate electrons from hadrons and the method used to evaluate the non-photon contributions from semi-leptonic decays of heavy flavor mesons. The observed nuclear modification factors, R_{AA} , of non-photon electrons indicate a substantial energy loss of heavy quarks at moderate to high p_T .

The measurement of heavy quark production helps expand our knowledge about the nuclear matter produced in heavy-ion collisions and furthers our understanding of the energy loss mechanism of quarks in the hot and dense medium. Due to the large mass of heavy quarks, suppression of small angle gluon radiation should reduce their energy loss and consequently, any suppression of heavy-quark mesons at high p_T is expected to be smaller than that observed for hadrons consisting of light quarks [1, 2, 3].

Typically, heavy quark meson production is studied by their reconstruction through hadronic decays, e.g. $D^0 \rightarrow K^- \pi^+$ [4]. An alternative way to infer information about heavy quark production in heavy ion collisions is the study of electrons from semi-leptonic decays of D and B mesons. This allows STAR to study charm and beauty production up to substantially larger p_T than is possible by exploiting the hadronic decay channels. There are several sources of measured electrons. We divide them into non-photon electrons (signal) and photon electrons (background). The non-photon electrons are mainly from semi-leptonic decays of heavy mesons and the Drell-Yan process. The background photon electrons are from γ conversions, and π^0 , η Dalitz and light vector meson decays.

The results presented in this paper were obtained from an analysis of data recorded with the STAR detector [5] in years 2003 (p+p, d+Au) and 2004 (Au+Au). The two main detector systems used in the analysis were the Time Projection Chamber (TPC) and the Barrel Electromagnetic Calorimeter (BEMC). For the data presented here, the BEMC was only half instrumented, limited to the range $0 < \eta < 1$ and $0 < \phi < 2\pi$, consisting of 2400 towers, each covering $(\Delta\eta, \Delta\phi)=(0.05, 0.05)$. A Shower Maximum Detector (SMD) is located approximately 5 radiation lengths inside each tower module. It allows the measurement of the shower shape in η and ϕ with high precision $(\Delta\eta, \Delta\phi)=(0.007,$

*For the full author list and acknowledgements see Appendix Collaborations in this volume

0.007). During the 2003 and 2004 runs, STAR used a high tower trigger (HT), based on the highest energy measured by a single tower in the BEMC, to enrich the event samples with high p_T particles that produce signal in the calorimeter, including the electrons. For the Au+Au data, we applied a single tower threshold of $E_T = 3$ GeV/c. For the results presented here, approximately 6.7 M minimum bias events, 2.6 M high tower trigger events and 4.2 M central events (10%) were used.

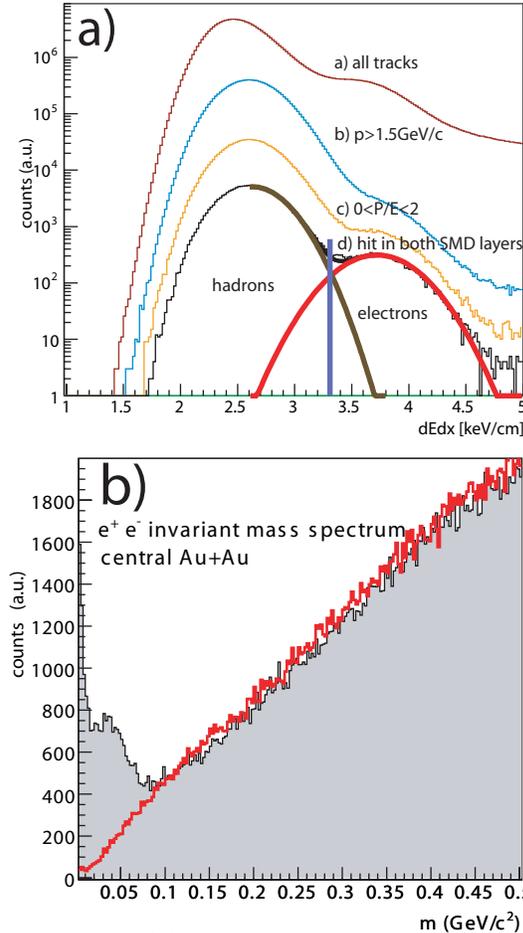


Figure 1. (a) TPC energy loss of tracks after applications of other electron criteria. (b) Invariant mass of reconstructed e^+e^- pairs (gray filled area) and combinatorial background (solid red line) in Au+Au central collisions.

A clean sample of electrons is obtained, with the p_T dependent residual hadron contamination that varies from 10 to 15%. The combined TPC+BEMC hadron discrimination power is p_T dependent and of the order of $10^2 - 10^4$.

The data sample for the Au+Au dataset was divided into 3 centrality bins (0-5%, 10-40%, and 40-80% most central). The electron reconstruction efficiency and acceptance were determined by embedding simulated electrons into real events and calculated for each centrality separately. For the most central events, the electron reconstruction efficiency increases with p_T up to 5 GeV/c and then remains constant at 40%.

The electron identification was based on a combination of energy loss in the TPC, energy deposition in the BEMC, and shower profile in the SMD. We shall refer to electrons in the text meaning both electrons and positrons. Only electron tracks reconstructed in the TPC which points to the interaction vertex are taken into account, thus eliminating a large fraction of conversion electrons. The measurement of the energy loss, dE/dx , in the TPC is a powerful method for particle identification. Hadrons with $p > 1.5$ GeV/c lose less energy than electrons with the same momentum and can be therefore well separated. The resulting electron candidates are then extrapolated to the BEMC towers. Their energy deposited in the calorimeter, E , is compared with the momentum, p , measured in the TPC. After traversing the TPC, the electrons deposit their total energy in the BEMC towers, yielding a p/E ratio close to unity, while hadrons showers are less developed and give $p/E > 1$. Electromagnetic showers are better developed than hadronic showers in the SMD detector, yielding to larger reconstructed clusters. Figure 1a shows the energy loss of the tracks after each electron identification cut. The resulting distribution is fit with two Gaussians to account for hadrons and electrons. Finally, by applying the dE/dx cut, the electrons can be well separated from hadrons. The fits to the dE/dx spectra allow us, in addition, to precisely determine the remaining hadron contamination.

The dominant sources of photonic electrons are γ conversions in the detector material, π^0 and η Dalitz decays. The contribution from other sources is negligible when compared to systematic uncertainties. The photonic background was evaluated by identifying both the electron and positron from conversion or Dalitz pair. The identified electrons were combined with all tracks of opposite charge that passed only loose electron dE/dx cuts to enhance pair reconstruction efficiency. The invariant mass spectrum of the reconstructed e^+e^- pairs is shown in Fig. 1b. The pairs sample contains true background pairs as well as random combinations (combinatorial background). The combinatorial background was estimated from like-sign pairs and is represented on the plot by the solid red line.

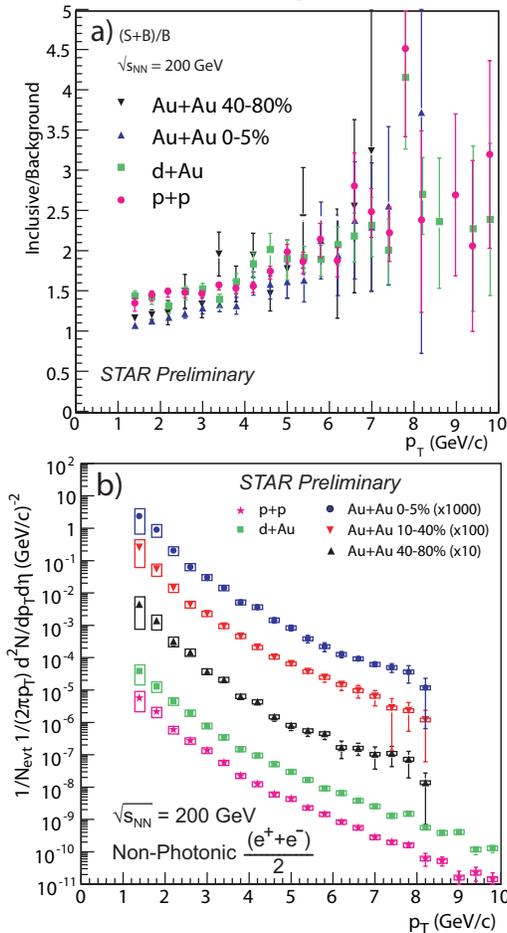


Figure 2. (a) Inclusive to photonic electrons ratio as a function of p_T for Au+Au collisions. (b) Background subtracted non-photonic electron spectra for p+p (pink), d+Au (green) and Au+Au collisions with centralities 0-5% blue, 10-40% red and 40-80% black.

The pairs from γ conversions and π^0 Dalitz decays have small invariant masses (see the counts over combinatorial background in Fig. 1b) and with a mass cut, $m < 140$ MeV/ c^2 , most of the photonic background was removed. The efficiency of the photonic background rejection was determined by embedding π^0 into real events and is about 60% for the most central Au+Au events and slightly decreases with p_T . The systematic uncertainties arising from the treatment of η Dalitz decays using π^0 decay kinematics was studied and found to be negligible for $p_T > 1.5$ GeV/ c . In Fig. 2a, the ratio of inclusive electrons to photonic background electrons is shown as a function of the p_T of the electrons. For $p_T > 2.0$ GeV/ c , there is a clear enhancement of electrons with respect to the background. This enhancement becomes more evident at higher momentum, where most electrons come from non-photonic sources, such as semi-leptonic decay of heavy quark mesons. Figure 2b shows the preliminary background subtracted non-photonic electron spectra for p+p, d+Au and Au+Au collisions. The error bars are statistical and the boxes show the preliminary systematic uncertainties. Many NLO pQCD, as well Pythia LO QCD calculations [6], predict that in a range between 3 and 6 GeV/ c the amount of electrons coming from B meson decays become significant. STAR is capable of measuring non-photonic electrons at a momentum range above this transition, making it possible to study in greater detail the interaction between heavy quarks and produced matter.

Figure 3 shows the nuclear modification factors R_{AA} and R_{dAu} for non-photonic electrons, as a function of p_T . The R_{dAu} ratio (Fig. 3a) is consistent with unity for the entire p_T range, suggesting that non-photonic electron production in d+Au follows a simple binary scaling with respect to p+p collisions. However, the errors are large and a mod-

erate Cronin type enhancement cannot be ruled out. On the other hand, it is possible to observe an increased suppression from peripheral to central Au+Au events (Fig. 3b-d) with respect to the binary scaling. Supposing that a significant fraction of non-photonic

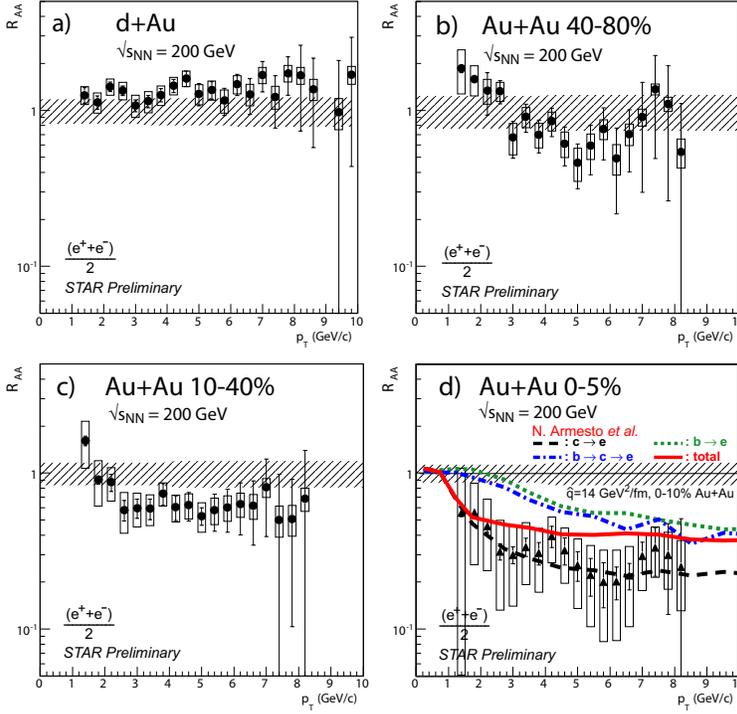


Figure 3. Nuclear modification factors R_{AA} for (a) identified hadrons at high p_T d+Au, (b) Au+Au 40-80%, (c) Au+Au 10-40%, and (d) Au+Au 0-5% comparison with model [3].

($R_{AA}(hadrons) \approx 0.2$ for central Au+Au), it suggests that the amount of energy loss for heavy quarks is comparable to the light quarks, raising questions about the understanding of the energy loss mechanism of partons in dense matter.

The non-photonic electron spectra measured by the STAR for p+p, d+Au, Au+Au collisions up to $p_T \approx 8$ GeV/c were presented. An increasing suppression of non-photonic electrons with the collision centrality in Au+Au collisions was observed. This may be related to a stronger than predicted interaction between heavy quarks and the medium created at RHIC. The analysis of the full statistics in the 2004 Au+Au run will permit a more detailed study of the medium modifications for heavy quarks allowing a better understanding of quark energy loss mechanisms. The higher statistics data set will also permit initial studies of correlations between non-photonic electrons with charged hadrons that will allow the study of heavy-quark tagged jets in Au+Au collisions at RHIC.

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